

High-energy hadron physics including possible spin projects at J-PARC

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Berkeley Summer Program on Nucleon Spin Physics

LBL, Berkeley, USA

<http://www-nsdth.lbl.gov/~fyuan/spin09/>

**Note: The latter half is based on my personal studies
rather than an overview of the J-PARC project.**

**June 1 – 12, 2009
(Talk on June 12)**

Contents

1. J-PARC facility

2. Low-energy hadron physics (→ short comments by skipping most slides)

Strangeness nuclear physics, Exotic hadrons,
Hadrons in nuclear medium, Neutrino-hadron interactions

3. High-energy hadron physics

- Hard interactions

Elastic scattering, Color transparency, Short-range NN interactions,
Generalized Parton Distributions (GPDs)

- Structure functions, Fragmentation functions

\bar{u} / \bar{d} asymmetry, Nuclear PDFs, Polarized PDFs,
Spin-1 structure, Fragmentation functions

- After major upgrades (polarization, heavy-ion ?)

Details of hadron spin, quark-hadron matters ...

3. Summary

J-PARC Facility

J-PARC Location

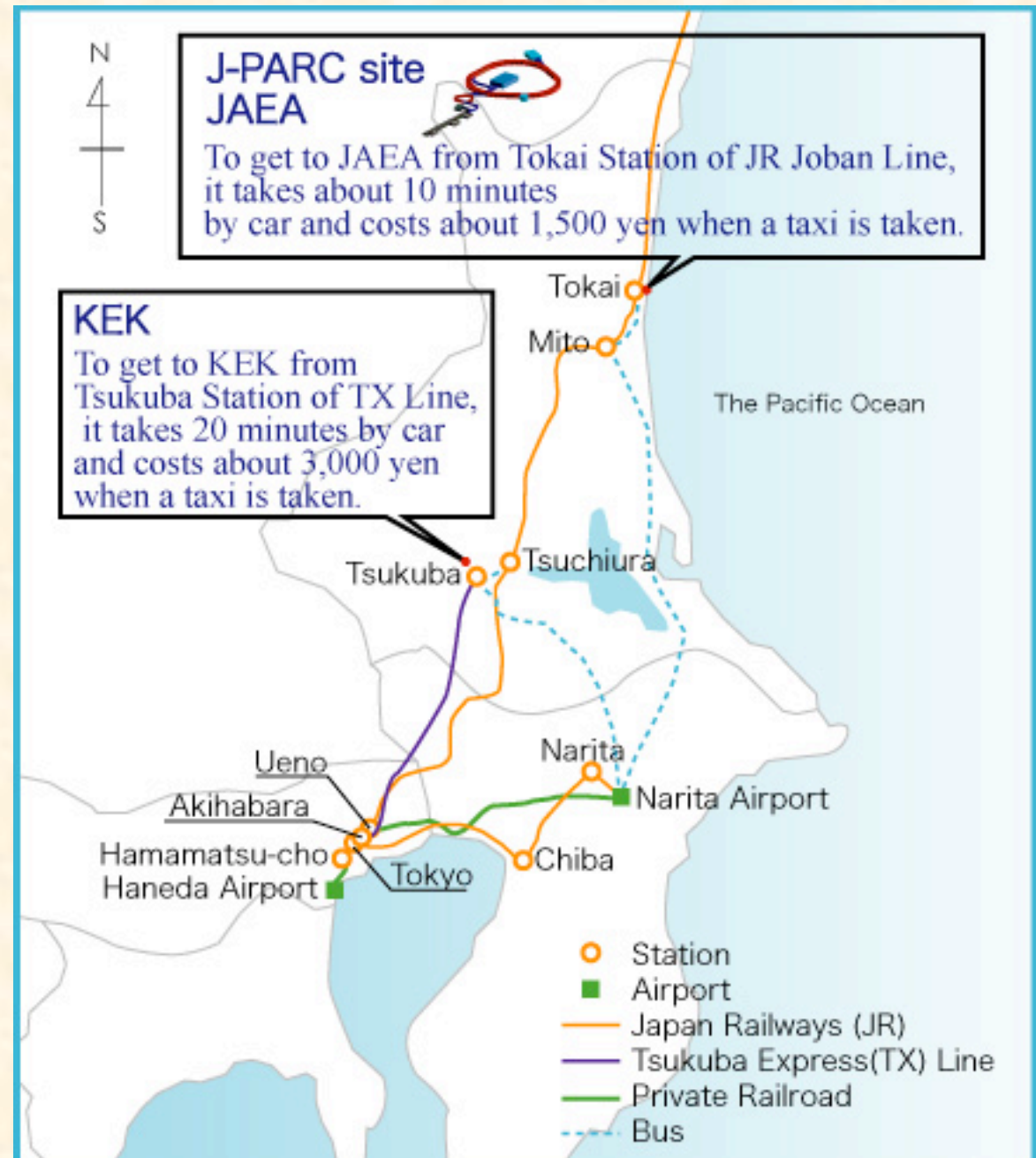
J-PARC (Japan Proton Accelerator Research Complex)

<http://j-parc.jp/index-e.html>

Joint facility of JAEA and KEK

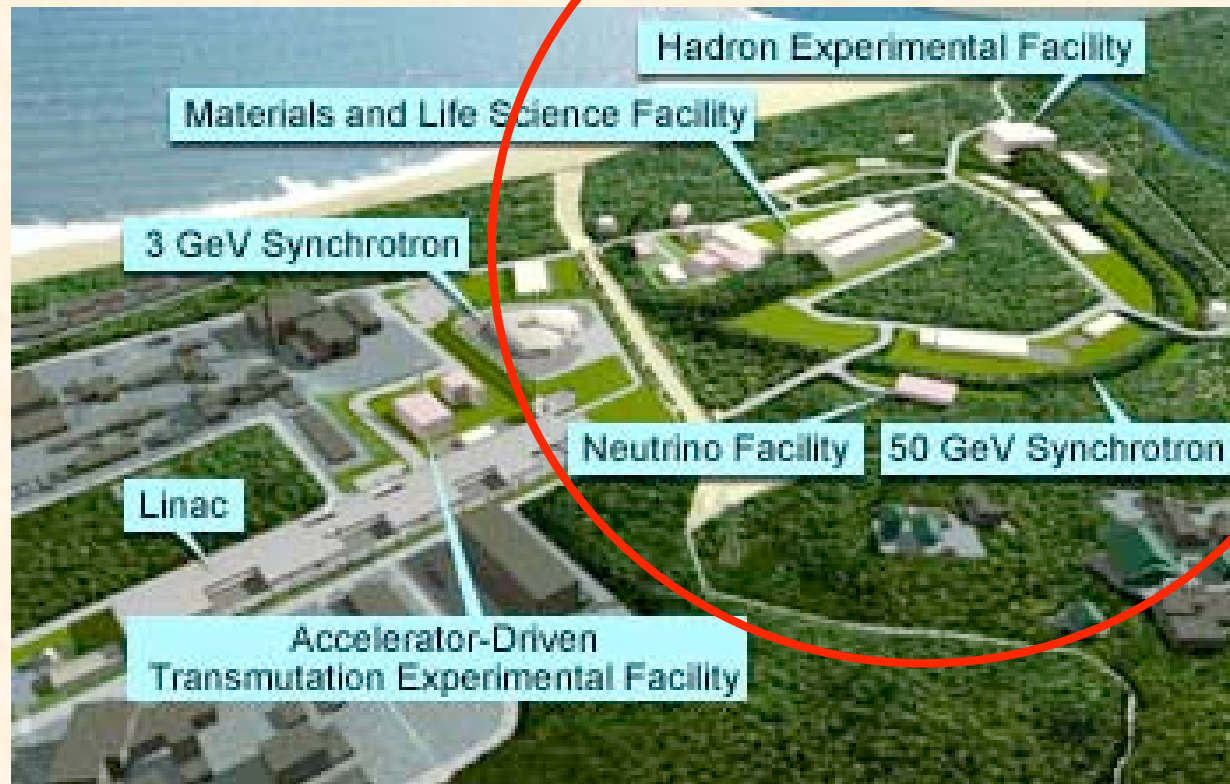
JAEA (Japan Atomic Energy Agency)

**KEK (High Energy Accelerator
Research Organization)**



Bird's-eye view

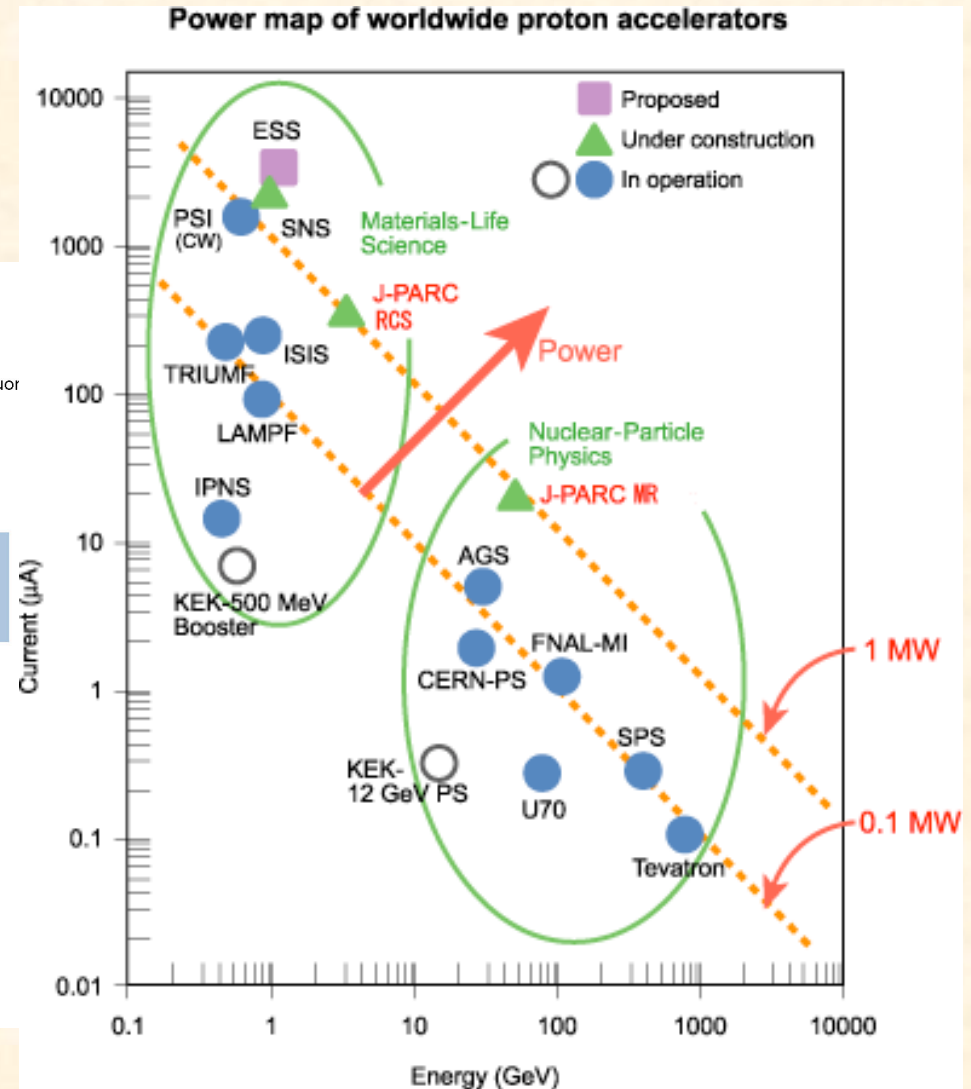
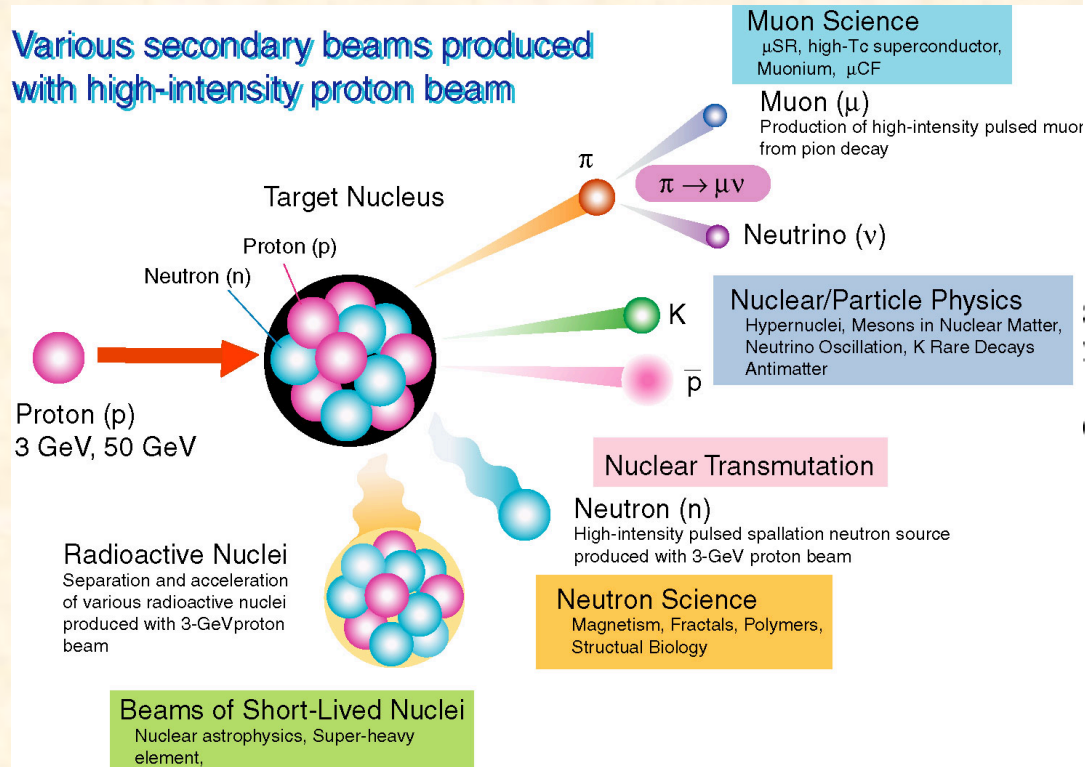
Particle and Nuclear Physics



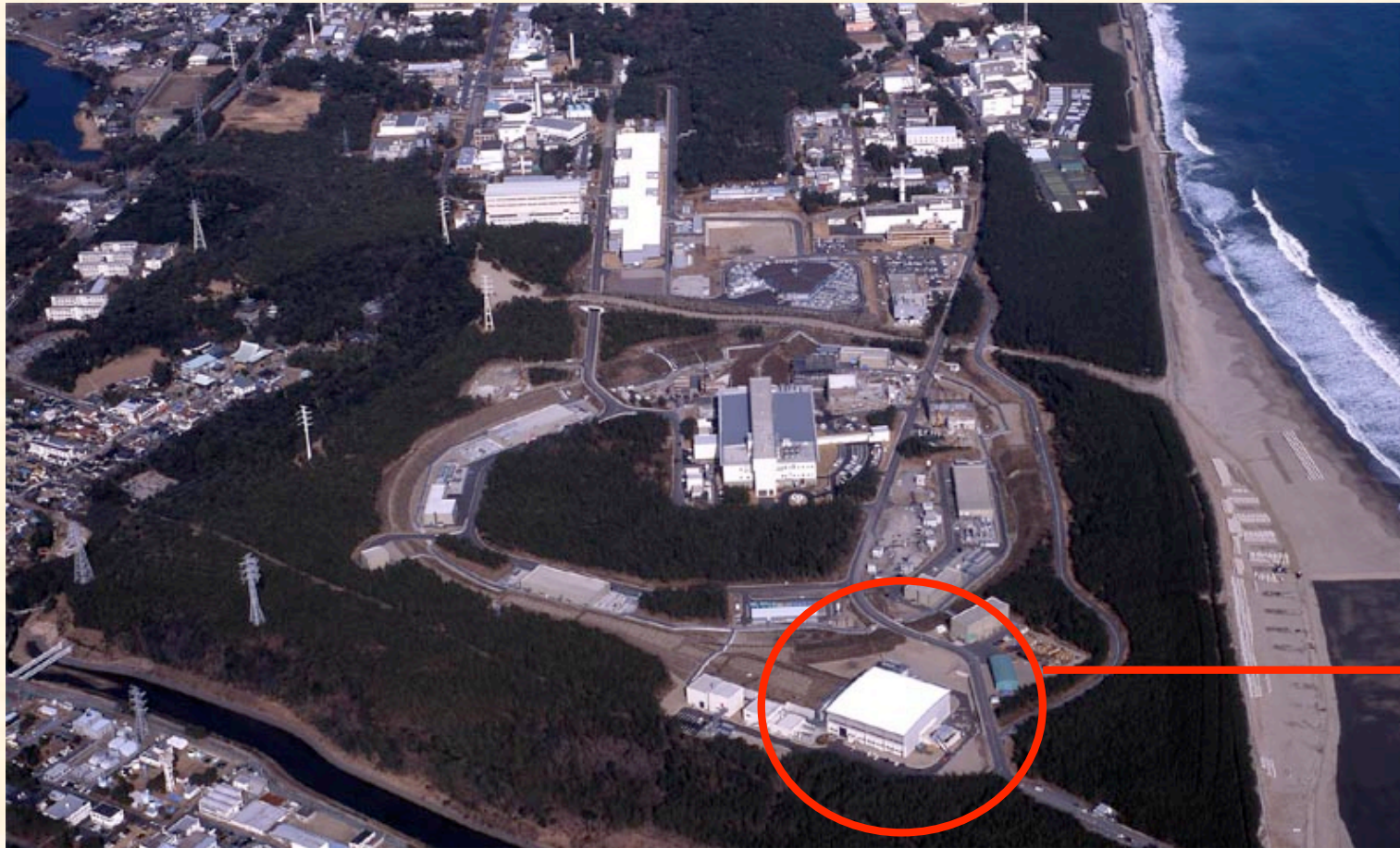
High-Intensity Frontier of Proton Accelerator

High-intensity proton beam
 → High-intensity secondary beams
 (Neutrino, Kaon, Pion, Neutron ...)

Various secondary beams produced
 with high-intensity proton beam



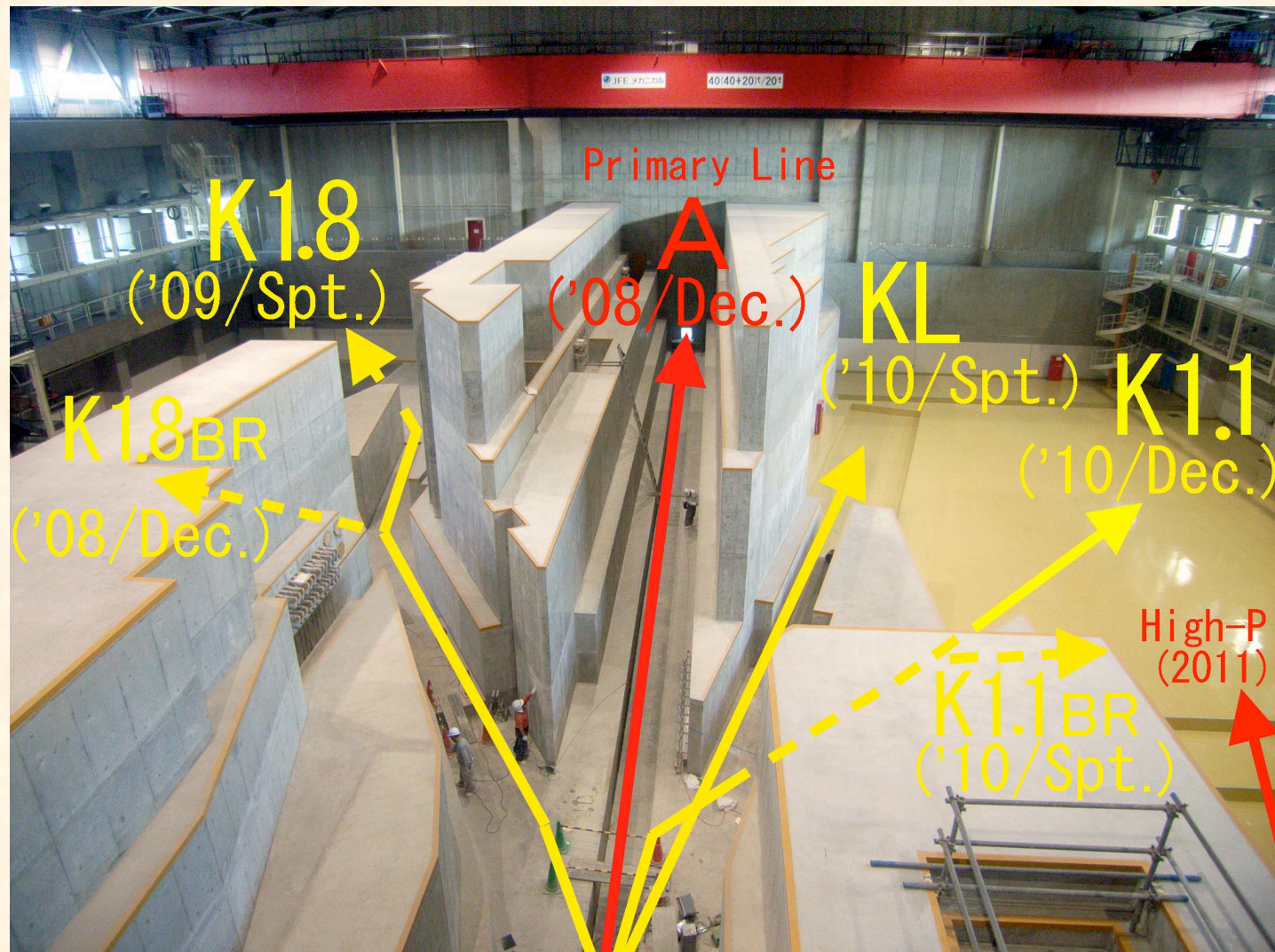
Aerial photograph on January 28, 2008



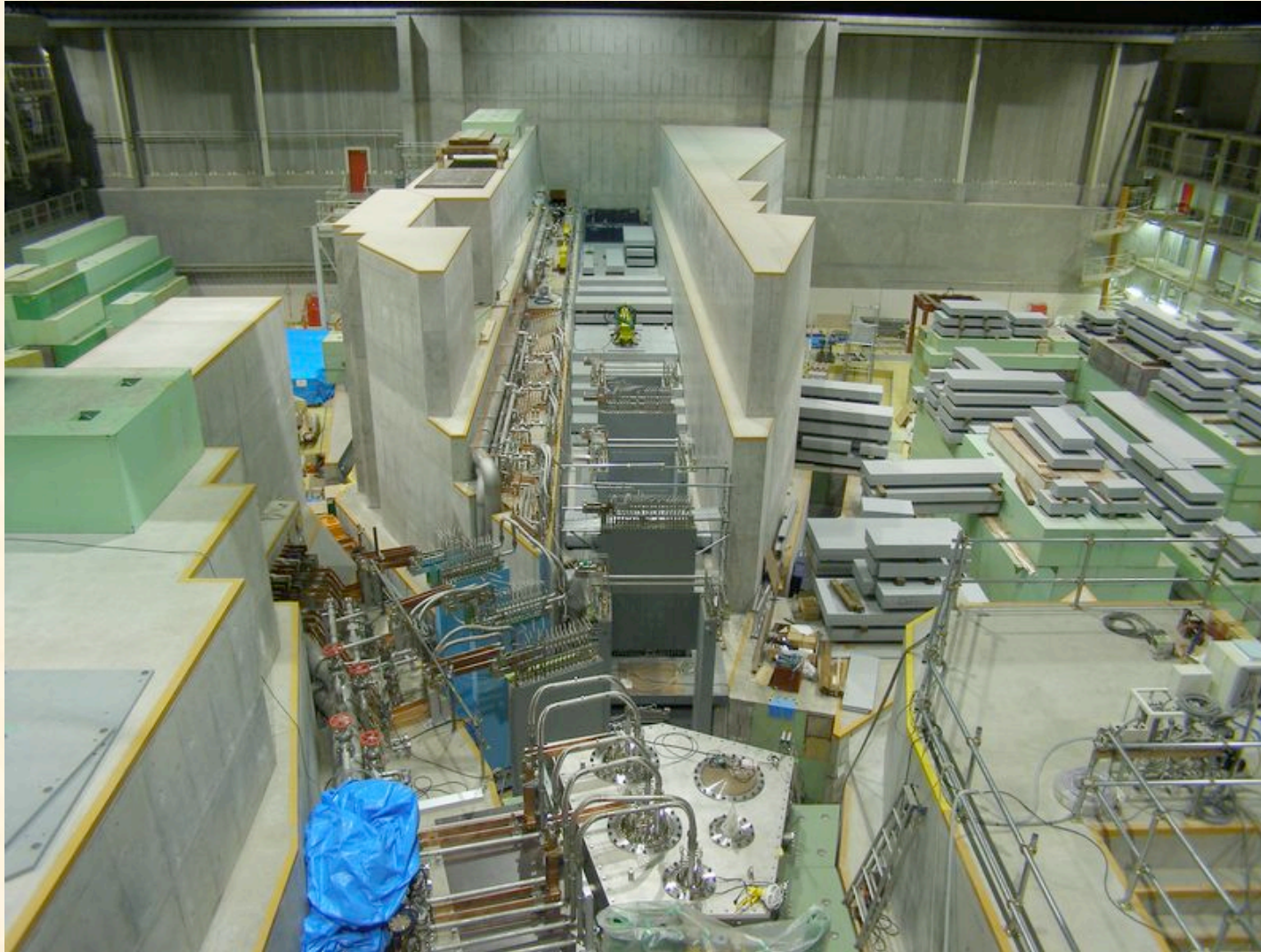
**Hadron
facility**

May 23, 2008: Injection to Main Ring
Jan. 28, 2009 : Acceleration to 30 GeV
Apr. 13, 2009 : Kaon beam
Apr. 23, 2009 : Neutrino beam

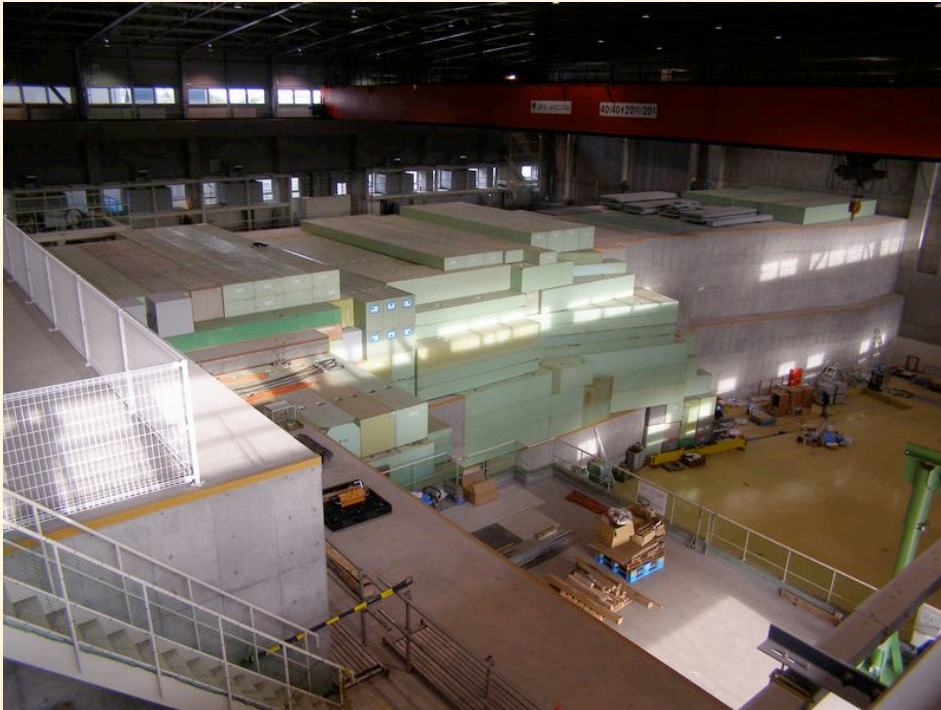
Hadron facility in May 2007 and a possible schedule for beam lines



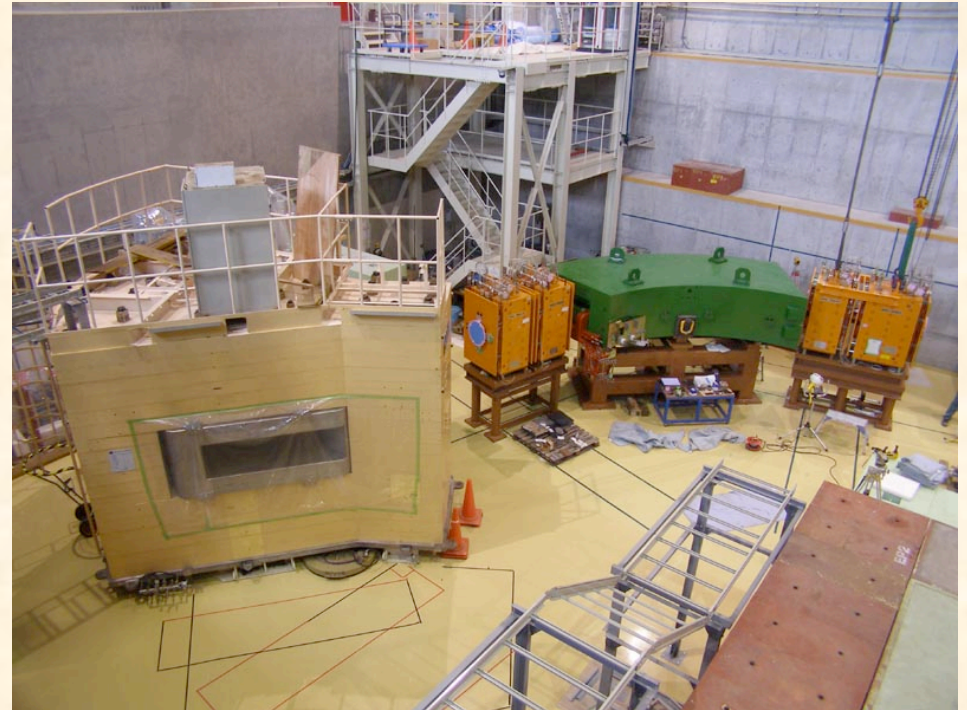
Hadron Facility on November 14, 2008



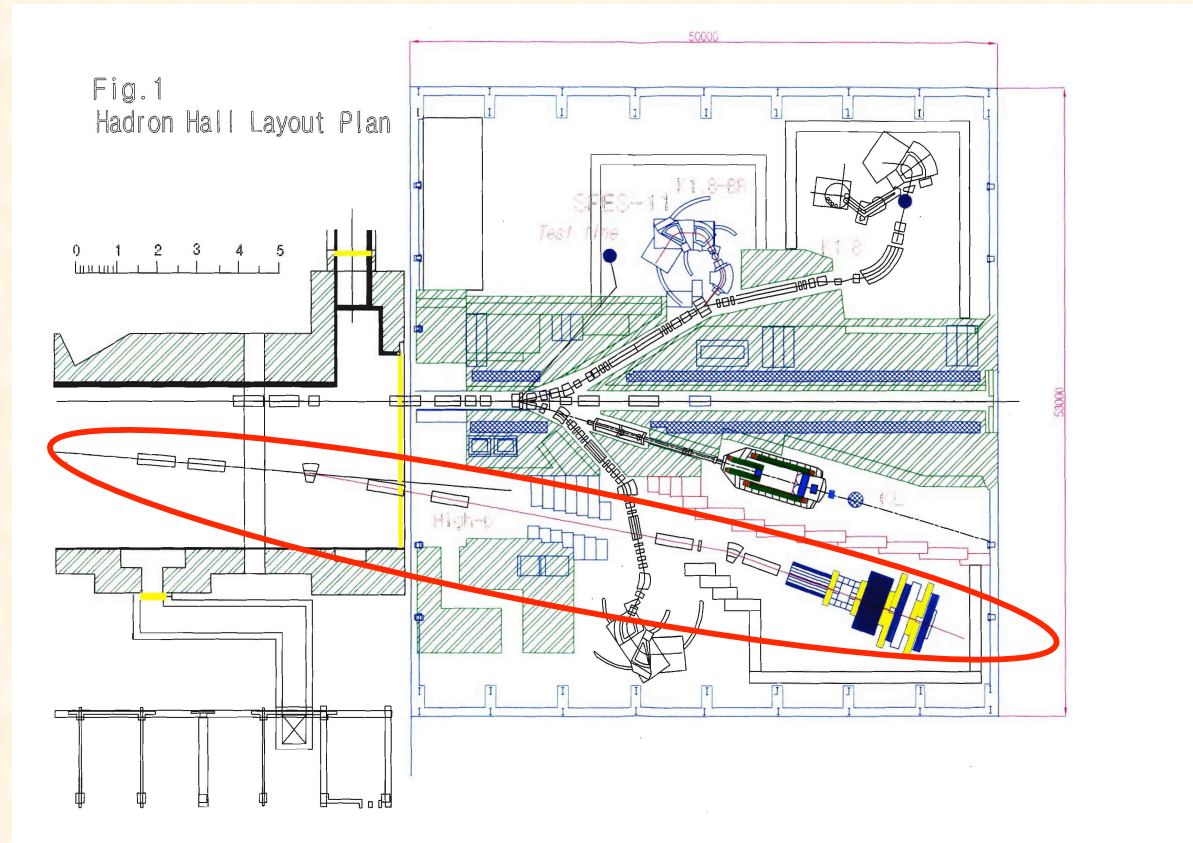
**Hadron Hall
on Jan. 13, 2009**



**K1.8 beamline
on May 12, 2009**



High-Momentum Beam Line (30, 50 GeV Proton)



This beam line should be interesting for the audience of this workshop.

Hadron Physics at J-PARC

by adding comments from my own works

References

My talks on “Possible Hadron Physics at J-PARC”

Trieste (2006) <http://www.pg.infn.it/hadronic06/>

Ghent (2007) <http://inwpent5.ugent.be/workshop07/>

Mito (2008) <http://j-parc.jp/NP08/>

Trieste (2008) <http://www.pg.infn.it/hadronic08/>

J-PARC workshops on hadron physics

- **J-PARC-HS05**

<http://www-conf.kek.jp/J-PARC-HS05/program.html>

- **J-PARC-NP07**

http://www-conf.kek.jp/NP_JPARC/program.html

- **J-PARC-NP08** <http://j-parc.jp/NP08/>

- **Riken spin workshop for J-PARC**

<http://rarfexp.riken.jp/~ygoto/jparc-riken0804/>

- **Meeting on high-energy hadron physics at J-PARC**

<http://www-conf.kek.jp/hadron08/hehp-jparc/>

J-PARC spin projects are discussed especially at these workshops.

Your contributions are welcome!
Please visit out KEK theory group.

Purposes of J-PARC hadron physics

**Understanding of strongly interacting matter
& Search for new state of matter**

Quantum Chromodynamics (QCD)

- **Asymptotic freedom, ...**
Perturbative QCD, Parton distribution functions,
Nucleon spin, ...
- **Color confinement**
Hadron spectroscopy,
Quark-hadron matter
- **Chiral symmetry**
Hadrons in nuclear medium

Hadron Physics at J-PARC

Efforts are needed to get approval for projects after strangeness physics.

1st project
(also ν)

- **Strangeness nuclear physics (1st experiment)**
Kaon and pion beams

- **Exotic hadrons**

- **Hadrons in nuclear medium**

Proton beam

- **Hard processes**

(50 GeV recovery)

- **Nucleon spin**

(proton polarization)

- **Quark-hadron matter (heavy ion)**

My talk is mainly on hard interactions.

Next projects

Need major upgrades

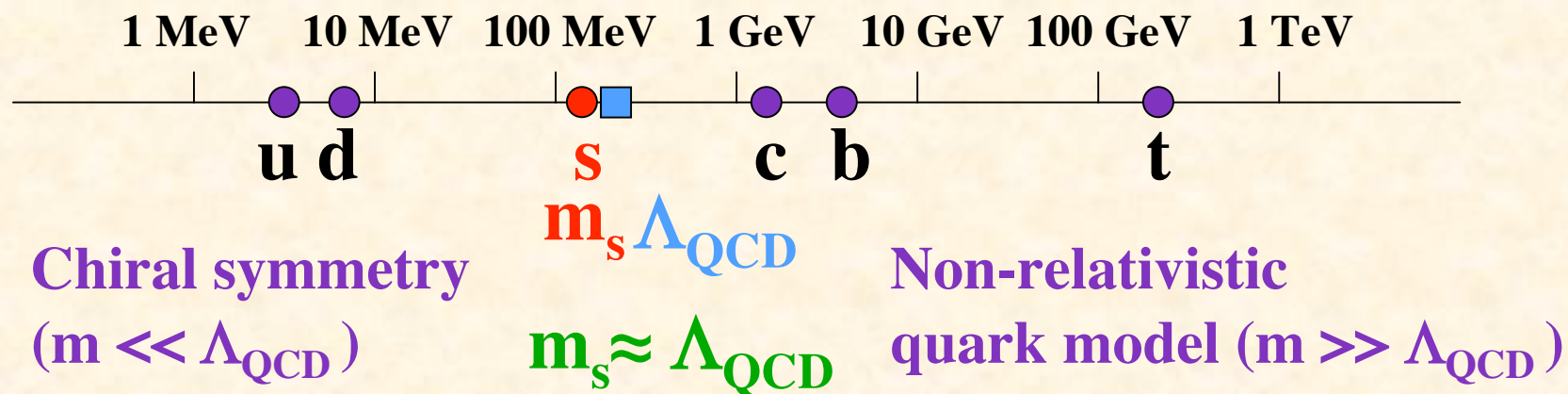
Low-Energy Hadron Physics

(I show you only a few slides on major projects.)

**I skip some transparencies on low-energy physics
because I should focus on hard-hadron interactions.**

Strangeness in hadron physics

(1) Probe of QCD dynamics



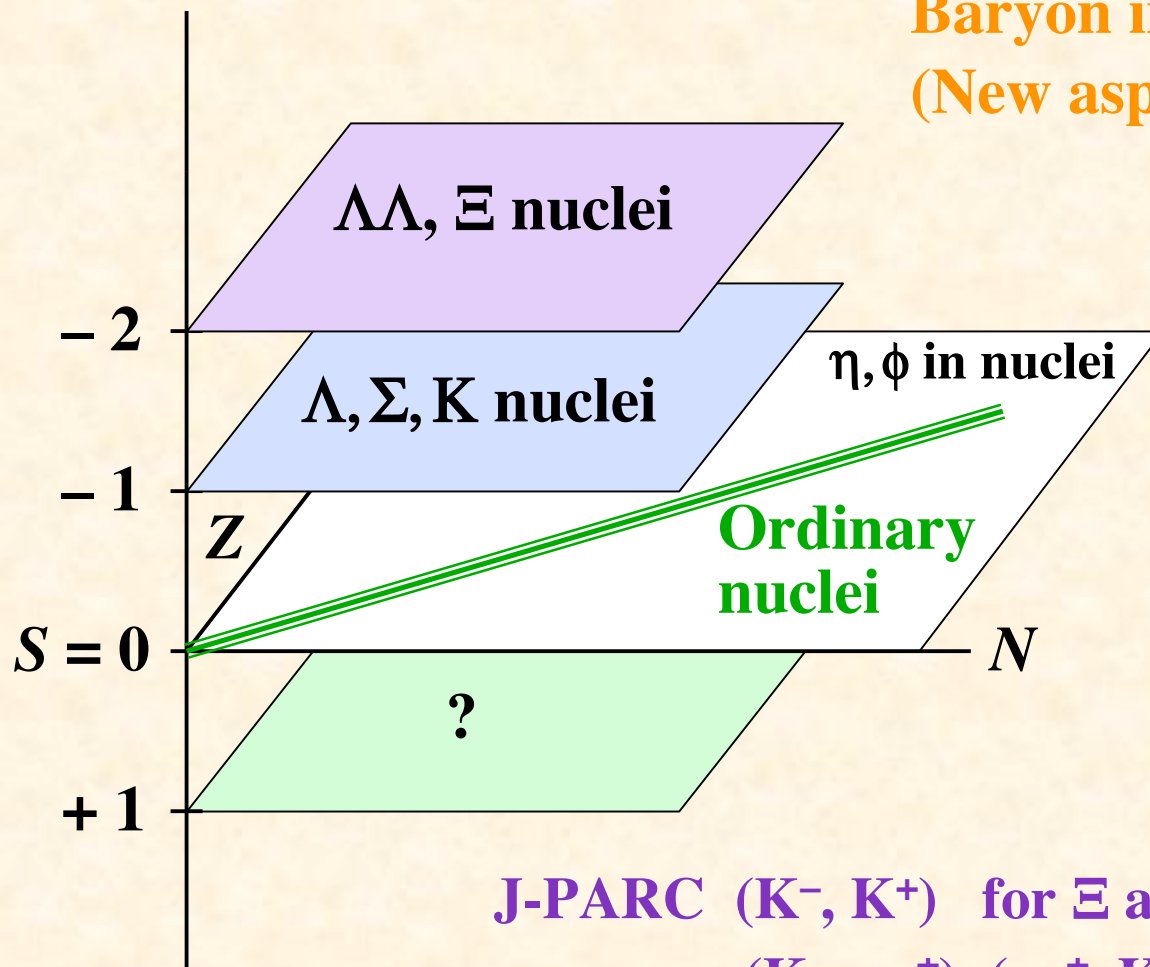
Bad point: It is difficult to describe hadrons with strangeness.

Good point: Strange quark could be a good probe of QCD dynamics.

(2) New nuclei with strangeness

New hadronic many-body system
by extending the flavor degrees of freedom.

**Baryon interactions with strangeness
(New aspect of low-energy QCD)**



- No data for YY interactions

- Some data for YN interactions (~ 40)

- Plenty of data for NN interactions ($\sim 4,000$)

J-PARC (K^- , K^+) for Ξ and $\Lambda\Lambda$ nuclei, YN scattering
(K^- , π^\pm), (π^\pm , K^+) for Λ nuclei, YN scattering

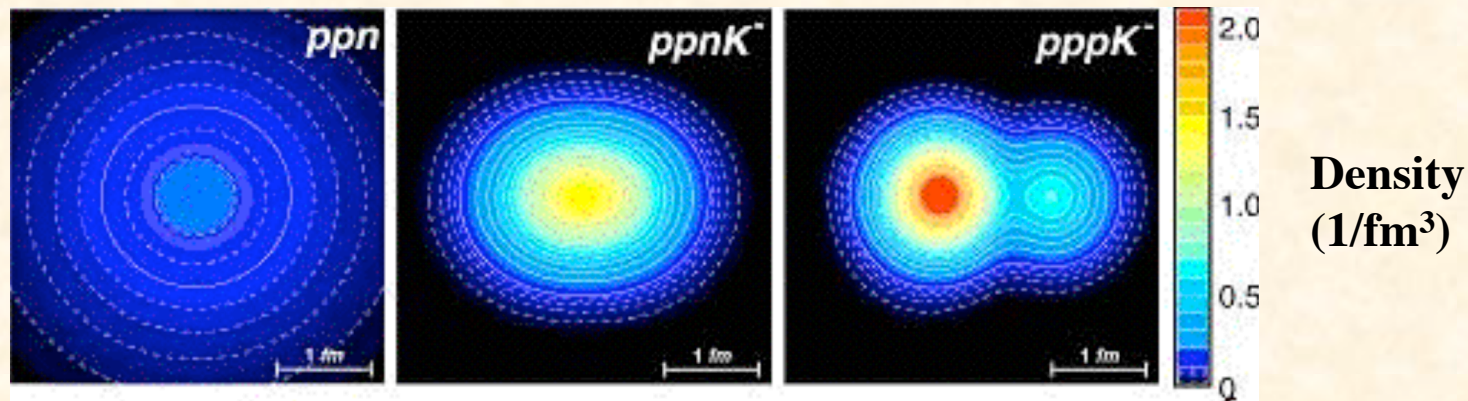
(3) Strangeness as impurity

No Pauli blocking with $u, d \rightarrow$ could penetrate deep inside a nucleus
(probe inside a nuclear medium)

New YN interactions \rightarrow could lead to new forms of nuclei
(possibly high-density nuclei)

Y. Akaishi, A. Dote, T. Yamazaki,
Phys. Lett. B613 (2005) 140.
See also Phys. Rev. C70 (2004) 044313.

- 1s atomic state of kaonic hydrogen
- $\bar{K}N$ scattering analysis
- Assume: $\Lambda(1405) =$ bound state of $\bar{K}N$
 \rightarrow Predictions of new kaonic nuclei



High-Energy Hadron Physics

Introduction:

Motivations and Issues

Comments on Structure Functions I

Advantages

- Studies of a different x region (large x)
from the ones at RHIC and LHC

(If the primary proton beam is polarized,
there is no rival facility in the world
for studying spin structure.)

- Different observables, *e.g.* antiquark and gluon,
from measurements at JLab (which is also a large- x facility)

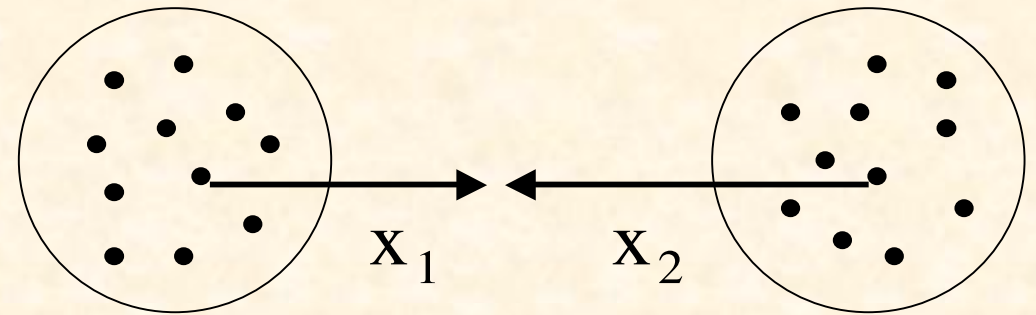
→ Next transparency for explanations

- ...

Hadron facilities

e.g. Drell-Yan: $x_1 x_2 = \frac{m_{\mu\mu}^2}{s}$

$\longrightarrow x \sim \frac{\sqrt{m_{\mu\mu}^2}}{\sqrt{s}}$



$$p + p(A) \rightarrow \mu^+ \mu^- + X \quad (q\bar{q} \rightarrow \mu^+ \mu^-)$$

- $s = (p_1 + p_2)^2$

J-PARC: $\sqrt{s} = 10 \text{ GeV}$

RHIC: $\sqrt{s} = 200 \text{ GeV}$

- $m_{\mu\mu} \geq 3 \text{ GeV}$

LHC: $\sqrt{s} = 14 \text{ TeV}$

$$x \sim \frac{\sqrt{m_{\mu\mu}^2}}{\sqrt{s}} \geq \frac{3}{10} = 0.3$$

J-PARC

**Large-x facility
(Medium-x)**

$$\geq \frac{3}{200} = 0.02$$

RHIC

$$\geq \frac{3}{14000} = 0.0002$$

LHC

Small-x facility

Comments on Structure Functions II

Advantage *and / or* disadvantage:

- **Perturbative QCD corrections are large.**
 - **Interesting for pQCD physicists**
 - **Can we obtain reliable parton distribution functions from measured cross sections?**
(→ next slides)

Applicability of perturbative QCD

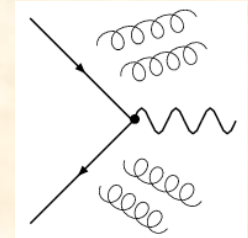
Cross section = pQCD \times non-pQCD (PDFs)

In order to extract the hadron-structure part,
pQCD should be understood.

Soft-gluon resummation is needed.

Large contributions come from the partonic threshold region

$$z = \frac{M_{\mu\mu}^2}{\hat{s}} \sim 1.$$



Drell-Yan cross section

Ref. H. Shimizu *et al.*, PRD 71 (2005) 114007

$$\frac{\tau d\sigma}{d\tau d\phi} \sim \sum_{a,b} \int_{\tau}^1 \frac{dx_a}{x_a} \int_{\tau/x_a}^1 \frac{dx_b}{x_b} f_a(x_a, \mu^2) f_b(x_b, \mu^2) \omega_{ab}(z, M_{\mu\mu}^2 / \mu^2, \alpha_s)$$

$$\tau = M_{\mu\mu}^2 / s, \quad z = \tau / (x_a x_b) = M_{\mu\mu}^2 / \hat{s}$$

e.g. in transverse spin asymmetry

$$\omega_{ab}(z, M_{\mu\mu}^2 / \mu^2, \alpha_s) = \omega_{q\bar{q}}^{(0)}(z) + \frac{\alpha_s}{\pi} \omega_{q\bar{q}}^{(1)}(z, M_{\mu\mu}^2 / \mu^2) + \dots$$

$$\omega_{q\bar{q}}^{(1)}(z, M_{\mu\mu}^2 / \mu^2) = C_F \left[4z \left(\frac{\ln(1-z)}{1-z} \right)_+ + \dots \right]$$

note: large contribution from the region $z \rightarrow 1$

Mellin transformation: $\int_0^1 dx x^{N-1} F(x)$

$$\frac{d\sigma^N}{d\phi} \sim \sum_f f^N(\mu^2) \bar{f}^N(\mu^2) \omega^N(M_{\mu\mu}^2 / \mu^2, \alpha_s)$$

$$\omega_{q\bar{q}}^{(1)N}(M_{\mu\mu}^2 / \mu^2) = C_F \left[2 \ln^2(N e^{\gamma_E}) + \dots \right]$$

A large term at $z \rightarrow 1$ corresponds to
a large term in the Mellin space at $N \rightarrow \infty$.

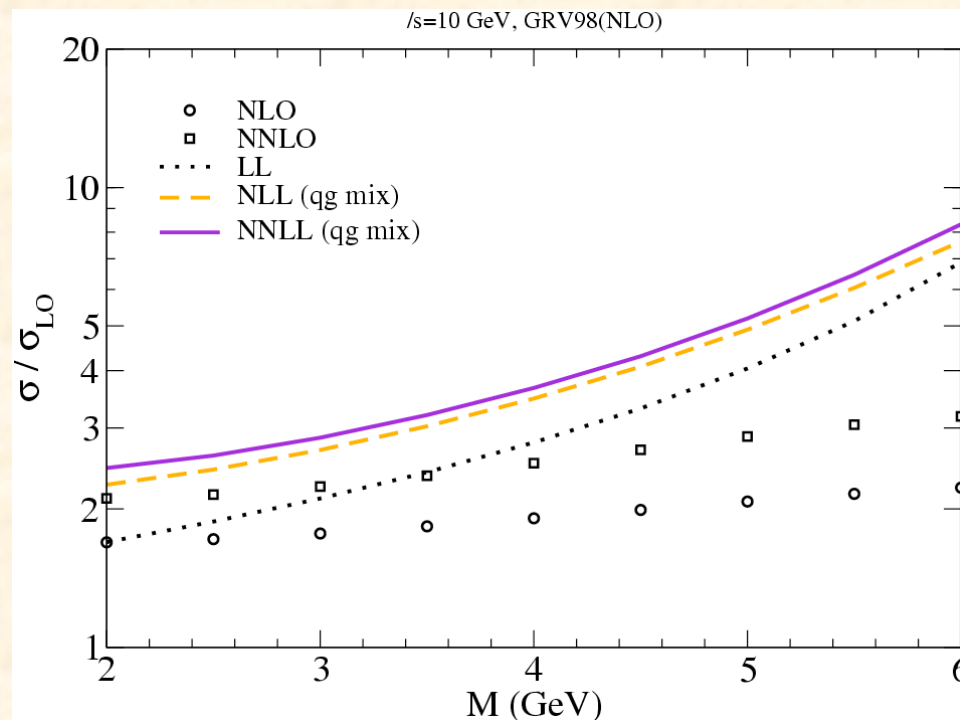
resummation

	fixed order	
	LO	1
NLO	$\alpha_s L^2$	$\alpha_s L$
\vdots	\vdots	\vdots
N ^k LO	$\alpha_s^k L^{2k}$	$\alpha_s^k L^{2k-1}$
	LL	NLL
	$L \equiv \ln N$	

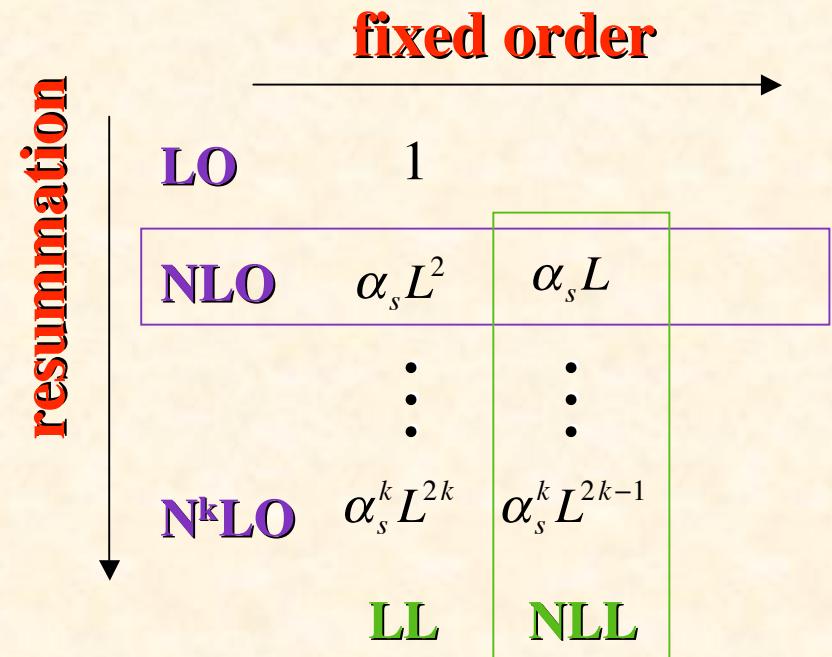
Applicability of perturbative QCD in Drell-Yan

- Higher-order α_s corrections
- Resummations

pQCD corrections are shown by $\frac{\sigma}{\sigma_{\text{Leading Order (LO)}}}$
as a function of the dimuon mass $M_{\mu^+\mu^-}$.



Yokoya@High-energy J-PARC
<http://www-conf.kek.jp/hadron08/hehp-jparc/>
 H. Yokoya and W. Vogelsang,
 AIP Conf. Proc. 915 (2007) 595.



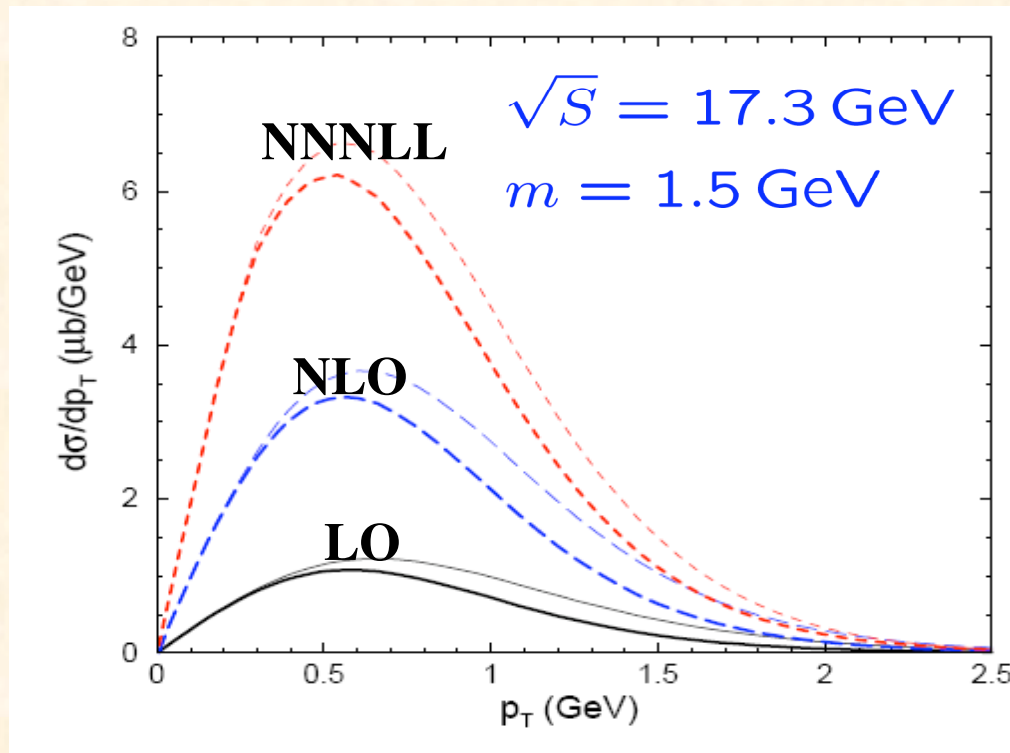
**Higher-order corrections are large at J-PARC (50 GeV);
 however, the pQCD terms could be under control in Drell-Yan.**

Applicability of perturbative QCD in charm production

N. Kidonakis, R. Vogt, Eur. Phys. J. C36 (2001) 201.

Stratmann@High-energy J-PARC

<http://www-conf.kek.jp/hadron08/hehp-jparc/>



J/ ψ production: see Z.-B. Kang, J.-W. Qiu, AIP Conf. Proc. 1056 (2008) 170.

Higher-order corrections are larger for charm-production processes, and their convergence is not studied as the level of the Drell-Yan case.

If someone can prove that theoretical calculations converge even at 30 GeV, high-energy projects (related to any PDFs) become appropriate at J-PARC. (Very Important !!!)

Comments on Structure Functions III

Motivations and Issues:

- **Why do you need such detailed nucleon structure?**
 - For example, the nucleon spin is one of fundamental physical quantities, but it is not understood.
In order to find its origin, we need measurements from small x (RHIC, (LHC)) to large x (JLab, J-PARC).
- **Costs for proton beamline, 50-GeV recovery, and proton-beam polarization.**

This is nothing to do with physics, but we should propose important experiments which are worth their costs.

 - ...

Comments on Structure Functions IV

Motivations and Issues:

- **Research purposes, Slogans for general public or at least for researchers in neighboring fields**

- **Applications ✓ Justifiable**

In calculating any hard interactions with a hadron, parton distribution functions are needed.

→ New physics (*e.g.* subquark, properties of quark-hadron matters, ...) at RHIC and LHC.

Possibly also to ultra-high energy cosmic rays ($\sim 10^{20}$ eV)

“High-energy nuclear physics at EeV (exa 10^{18}) and ZeV (zetta 10^{21}) ”

- **Fundamentals**

Nucleon spin, Confinement mechanism, ...

- **Appealing slogan(s)?**

... (Ideas of young people?)

Quark substructure?

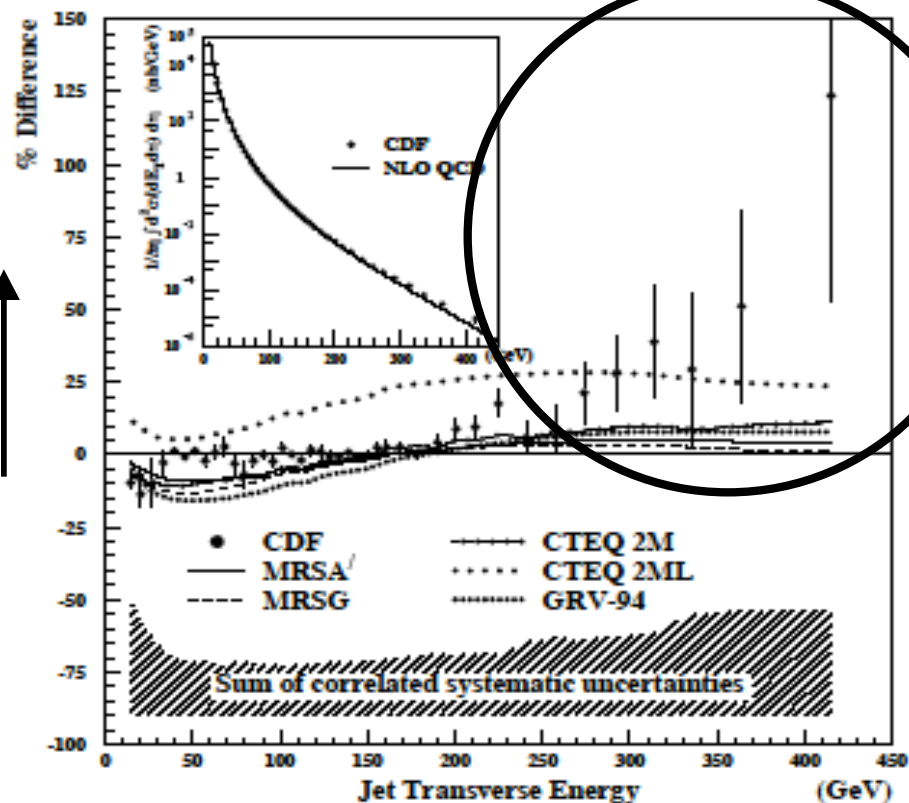
CDF experiment: PRL, 77 (1996) 438.

Comparison of theoretical calculations
with CDF experimental data.

$$p + \bar{p} \rightarrow \text{jet} + X$$

$$\sqrt{s} = 1.8 \text{ TeV}, \quad E_T^{\text{jet}} = 15 - 400 \text{ GeV}$$

Difference between theory and experiment



Jet transverse energy

Subquark signature ???

The same thing could happen
at LHC.

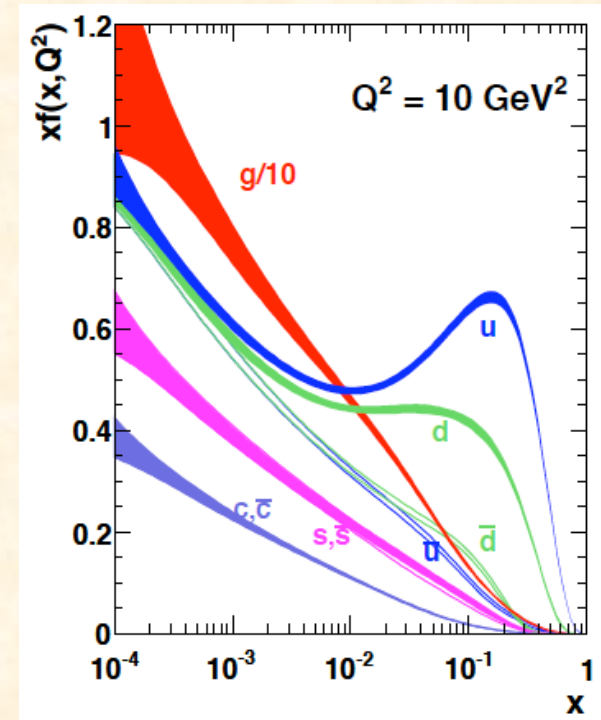
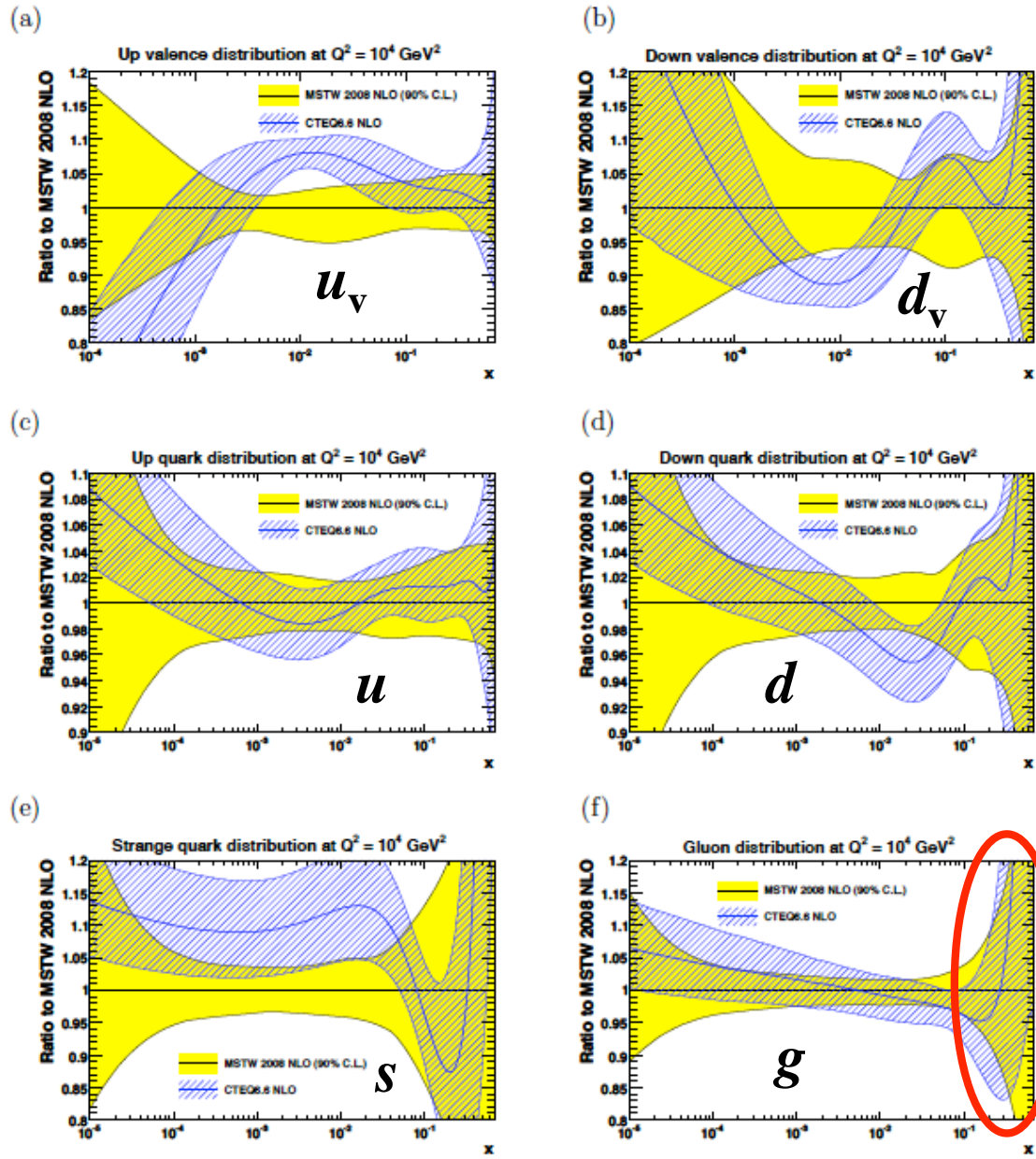
Could be explained
without substructure

(importance of accurate PDFs)

PDF uncertainty

A. D. Martin, W. J. Stirling, R. S. Thorne,
and G. Watt, arXiv: 0901.0002.

 MSTW
 CTEQ6.6



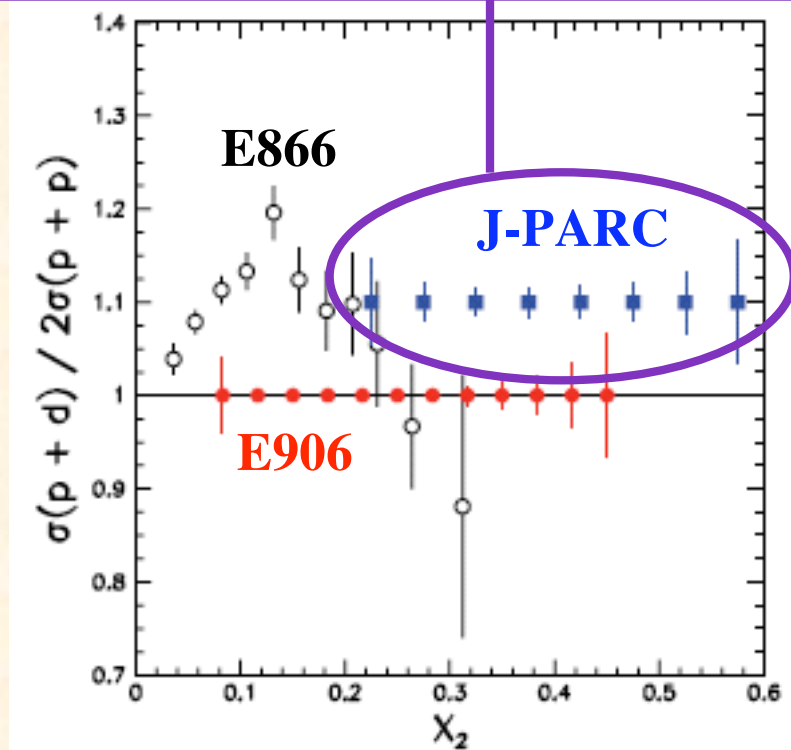
Important x region for finding
an “exotic event” in a high- p_T
region at LHC.

J-PARC x region

Possible Topics I

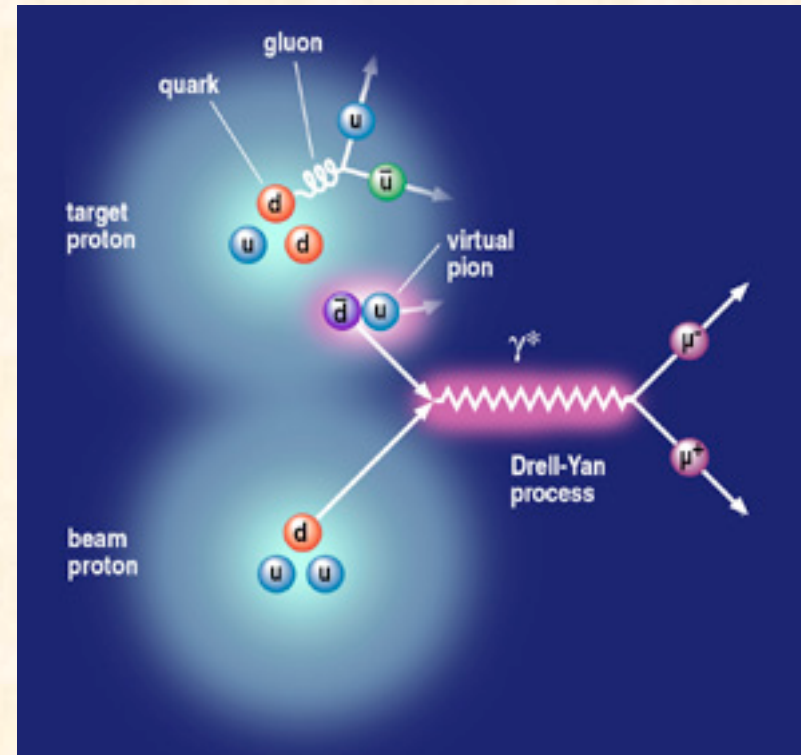
Flavor asymmetric antiquark distributions: \bar{u} / \bar{d}

Theoretical studies are needed for physics importance in this x region.



J-PARC proposal (P24), M. Bai *et al.* (2007)

This project is suitable for probing
“peripheral structure” of the nucleon.



<http://www.acuonline.edu/academics/cas/physics/research/e906.html>

Refs. SK, Phys. Rep. 303 (1998) 183;
G. T. Garvey and J.-C. Peng,
Prog. Part. Nucl. Phys. 47 (2001) 203.

Spin asymmetry in elastic scattering

Single spin asymmetry

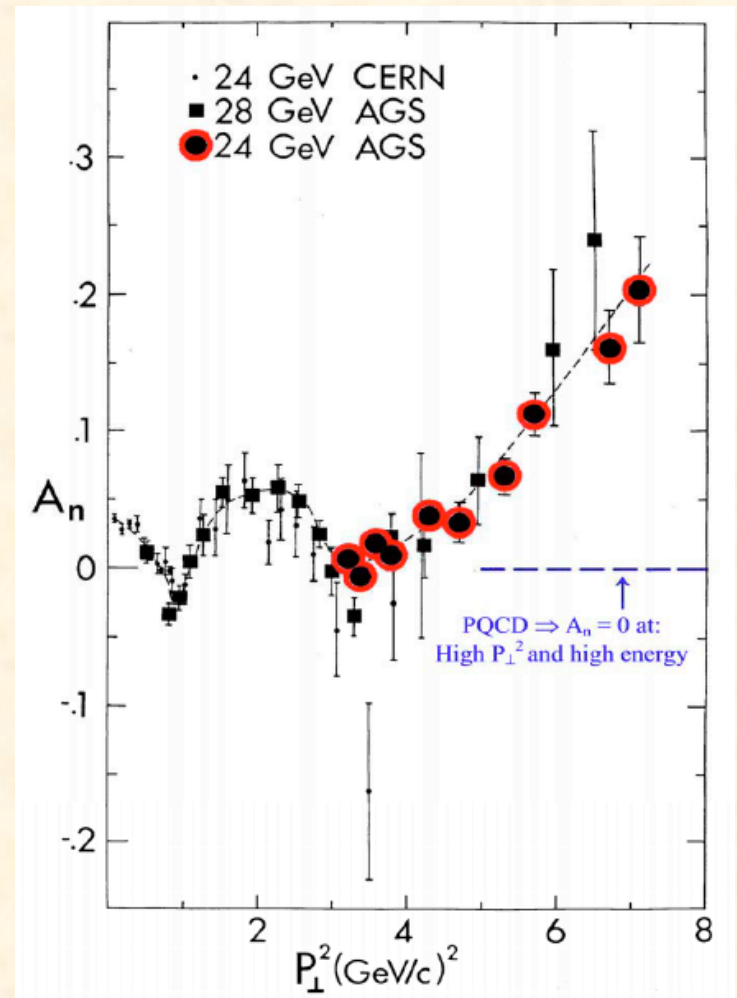
$$A_n = \frac{\sigma^{\uparrow} - \sigma^{\downarrow}}{\sigma^{\uparrow} + \sigma^{\downarrow}}$$

- J-PARC 30 GeV is the same as the AGS energy.
(Kinematical range is similar.)

Simply the AGS type measurements do not have physics impact.

For a possible J-PARC experiment,

- New observable should be investigated for providing a clue to pin down a possible mechanism of producing the asymmetry at large p_T .
- Suggestions from theorists?



Krisch@Riken-J-PARC08

D. G. Crabb et al., PRL65 (1990) 3241

Possibilities are discussed in the proposal P24 (Y. Goto *et al.*)

Proposal P24: Polarized proton acceleration at J-PARC

Proposal

Polarized Proton Acceleration at J-PARC

November 30, 2007

M. Bai¹, M. Brooks⁵, J. Chiba¹¹, N. Doshita¹², Y. Fukao⁷,
Y. Goto^{7,8†}, M. Grosse Perdekamp², K. Hatanaka⁶, H. Huang¹,
K. Imai⁴, T. Iwata¹², S. Ishimoto³, X. Jiang⁵, K. Kondo¹²,
G. Kunde⁵, K. Kurita⁹, M. J. Leitch⁵, M. X. Liu⁵, A. U. Luccio¹,
P. L. McGaughey⁵, A. Molodjontsev³, C. Ohmori³, J.-C. Peng²,
T. Roser¹, N. Saito³, H. Sato^{3†}, S. Sawada³, R. Seidl²,
T.-A. Shibata¹⁰, J. Takano³, A. Taketani^{7,8}, M. Togawa⁸, and
A. Zelenski¹

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Wako, Saitama 351-0198, Japan

⁸ RIKEN BNL Research Center, Brookhaven National Laboratory,
Upton, NY 11973, USA

⁹ Rikkyo University, Tokyo 171-8501, Japan

¹⁰ Tokyo Institute of Technology, Tokyo 152-8551, Japan

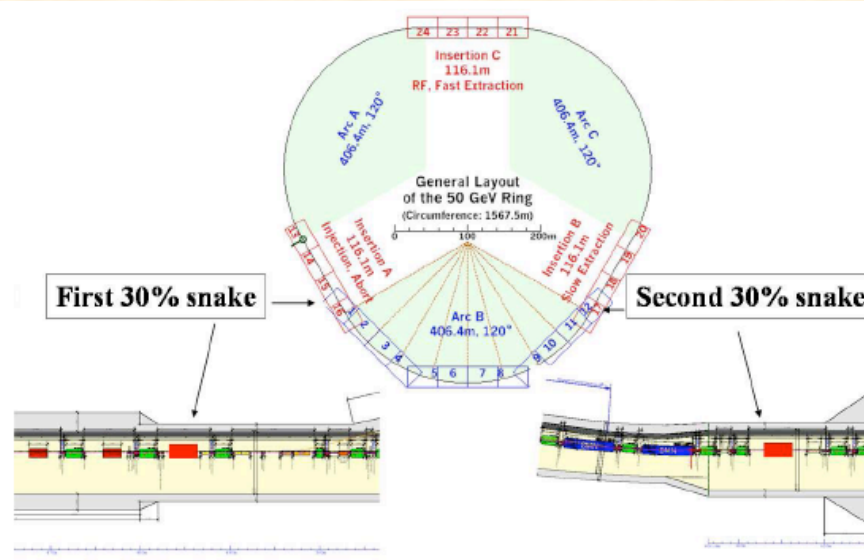
¹¹ Tokyo University of Science, Noda, Chiba 278-8510, Japan

¹² Yamagata University, Yamagata 990-8560, Japan

I introduce some topics from this proposal,
but please look at this proposal for technical
details of the beam polarization and physics
topics.

Primary-proton-beam polarization

Proposal P24: Polarized proton acceleration at J-PARC



The proton beam polarization is technically possible.

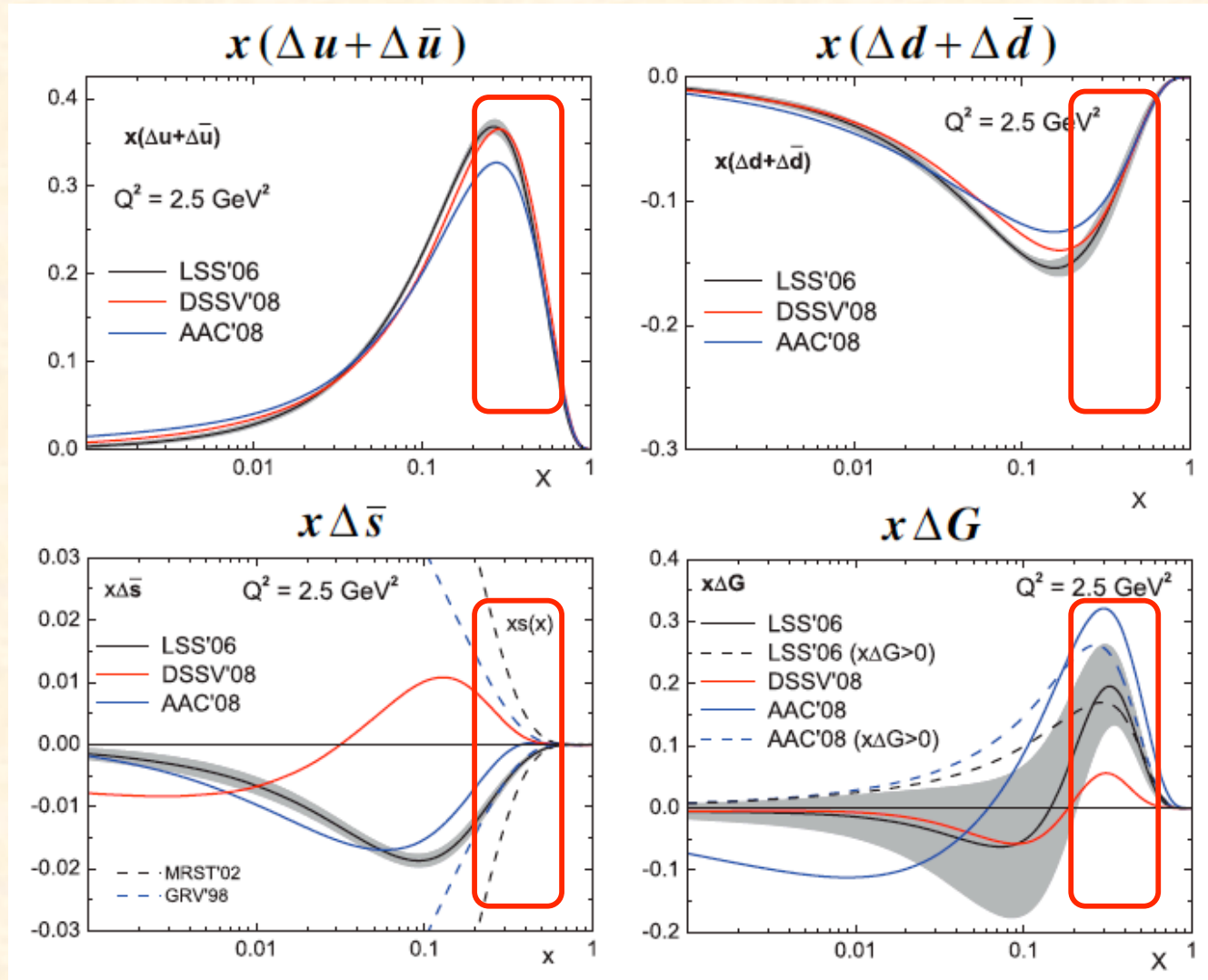
Figure 16: Possible location of partial helical snake magnets in the MR.

Item	Cost (million yen)
Polarized ion source	200
Source to RCS	50
Polarization in RCS	100
Polarized beam in MR	500
Primary beam extraction	250
pC polarimeter	100
Absolute polarimeter	100 - 300
Total	1,300 - 1,500

Table 1: Rough cost estimation for the polarized proton acceleration at J-PARC.

Situation of polarized PDFs

S. E. Kuhn, J.-P. Chen, and E. Leader,
arXiv:0812.3535 [hep-ph],
to be published in Prog. Part. Nucl. Phys.



**J-PARC
kinematical
range**

Gluon and antiquark distributions have large uncertainties.

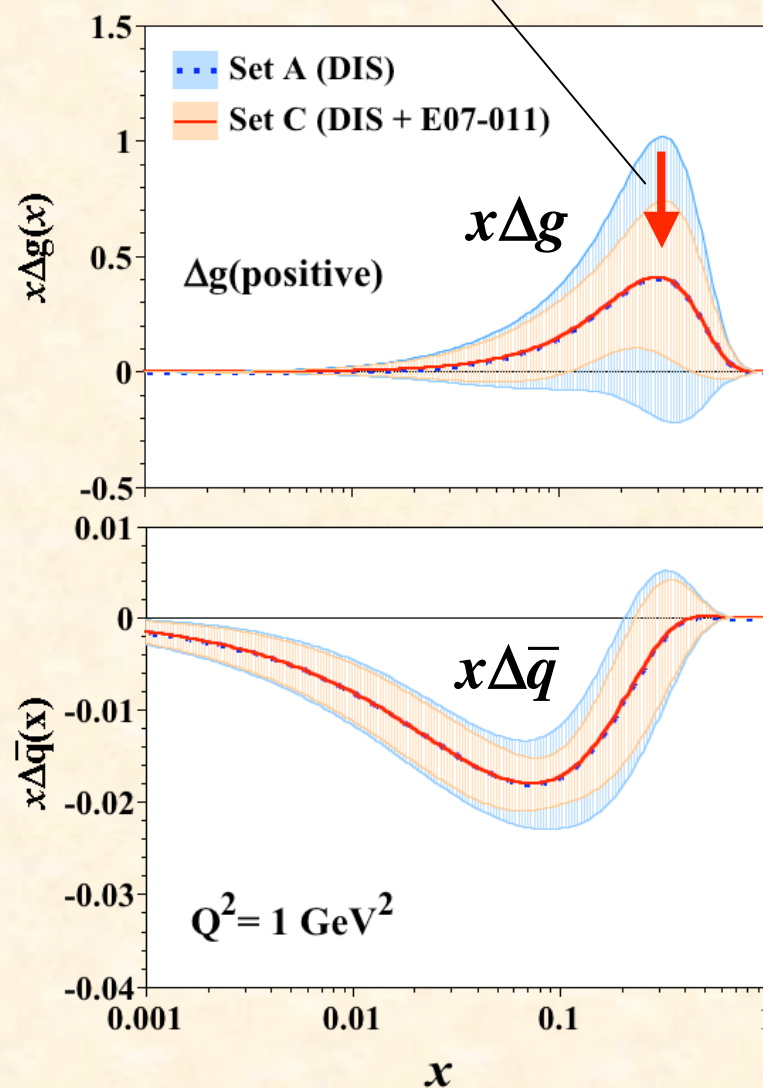
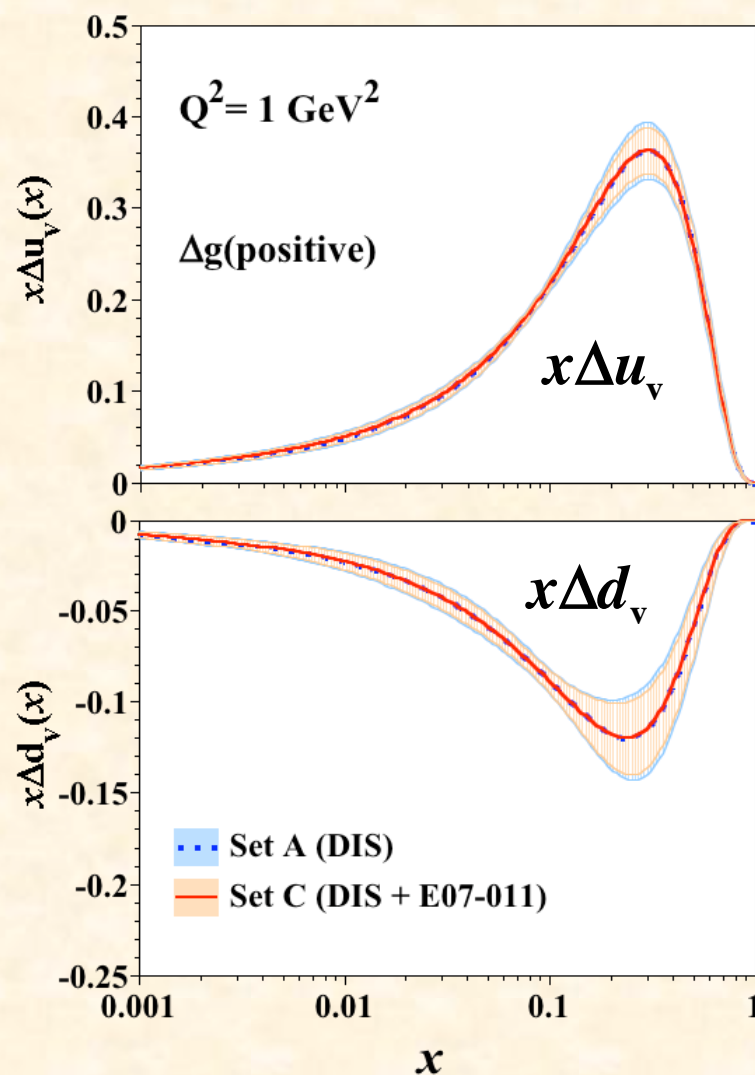
Comments on AAC08 polarized PDFs

**Determination of gluon polarization
from deep inelastic scattering and collider data,
M. Hirai and S. Kumano,
Nucl. Phys. B813 (2009) 106.**

Effects of future JLab E07-011 data

It is interesting to find the reduction of Δg uncertainties by the lepton scattering (JLab E07-011) data.

The same effect is expected at future EIC.



A: Current DIS data
C: Current DIS data + E07-011

Reason for the reduction of Δg uncertainty

Two possibilities:

- (1) Error correlation between antiquark and gluon distributions
(but note antiquark reduction is smaller).
- (2) Gluonic NLO term in g_1 (left figures).

$$\frac{1}{g_1(x, Q^2)} \frac{1}{2} \sum_{i=1}^{n_f} e_i^2 \int_x^1 \frac{dz}{z} \Delta C_g(x/z, Q^2) \Delta g(z, Q^2), \quad (5)$$

CLAS data are not accurate enough to probe the gluonic NLO term,

but **E07-011 data are so accurate that**

Δg is restricted through the NLO term in g_1 .

Q^2 dependence is not the reason for the reduction according to the figure below.

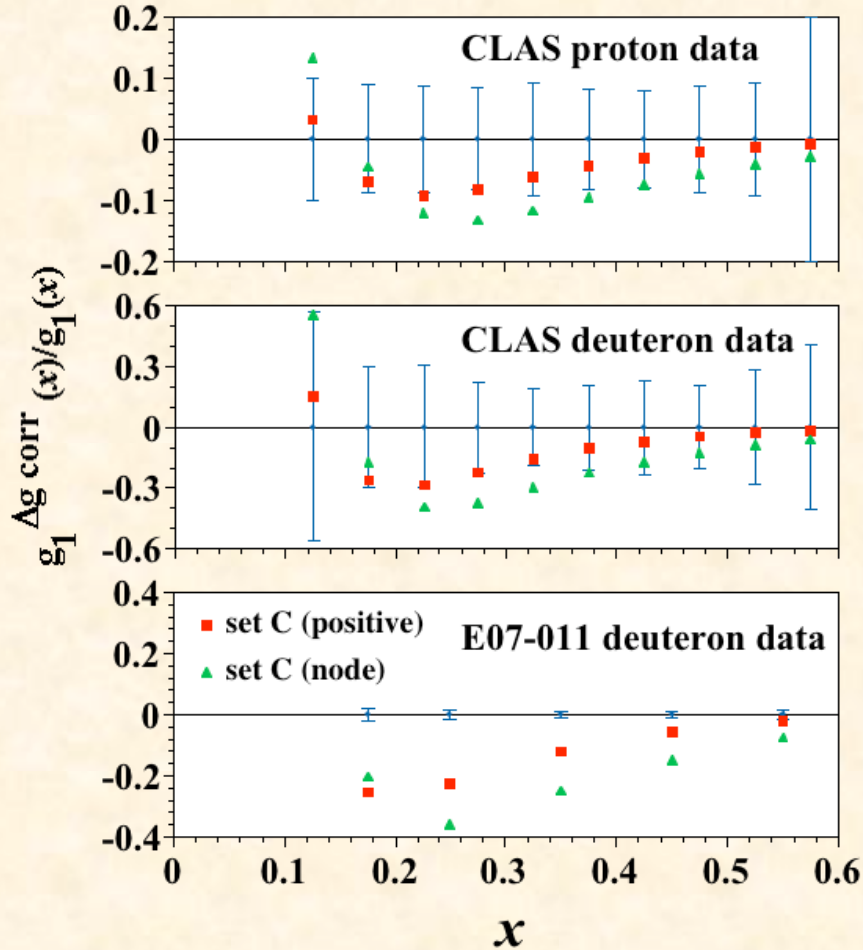
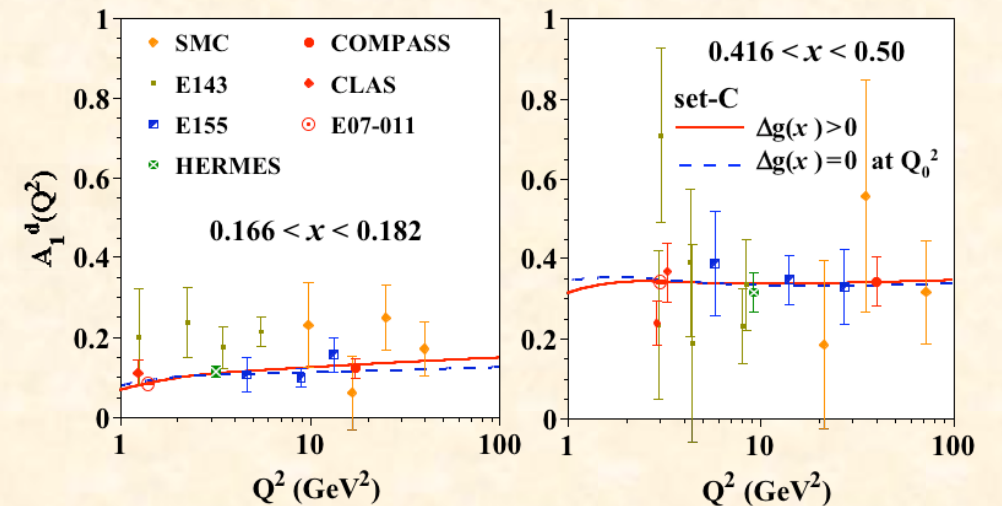
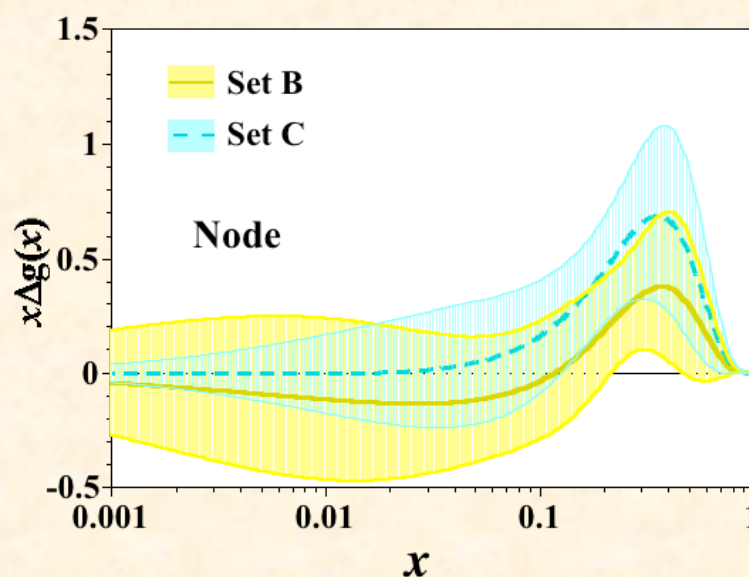
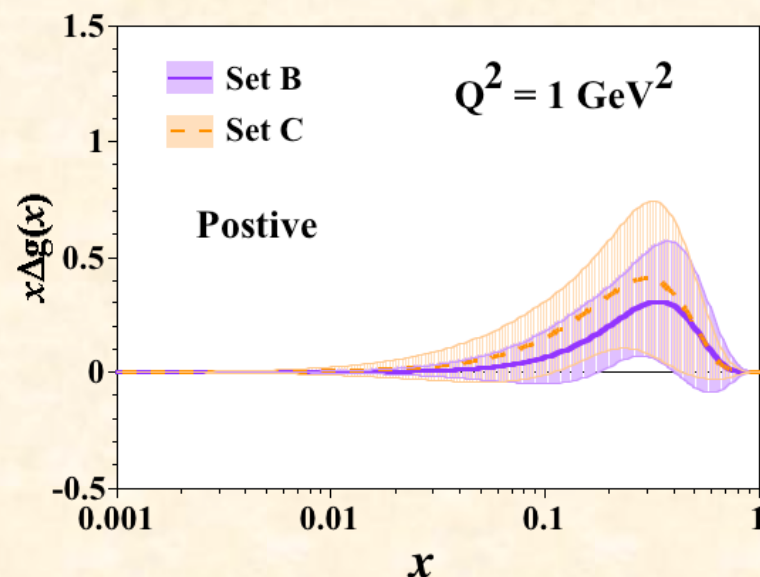


FIG. 8: The ratio of the gluon NLO-correction term to the polarized structure function $g_1(x, Q^2)$ in Eq. (5). It is compared with experimental errors, which are shown by the ratio $\delta g_1(x, Q^2)/g_1(x, Q^2)$ in Eq. (6).



$\Delta g(x)$ with PHENIX run-5 or JLab E07-011 data



B: Current DIS data
+ RHIC- π^0 (run-5)
C: Current DIS data
+ E07-011

Δg function	First moment	DIS	DIS+RHIC π	DIS+E07-011
positive	Δg	0.53	0.36	0.53
positive	$\delta(\Delta g)$	0.72	0.26	0.38
positive	$\delta(\Delta g)/\Delta g$	1.36	0.71	0.73
Significant improvements				
node	Δg	0.87	0.4	0.87
node	$\delta(\Delta g)$	0.89	0.31	0.47
node	$\delta(\Delta g)/\Delta g$	1.02	0.77	0.54

JLab-E07-011 is comparable to RHIC run-5 π^0 in determining $\Delta g(x)$.

Return to possible J-PARC projects

Single spin asymmetry (No polarized proton beam is needed!)

$$A_N = \frac{\sigma^\uparrow - \sigma^\downarrow}{\sigma^\uparrow + \sigma^\downarrow}$$

- Sivers effect

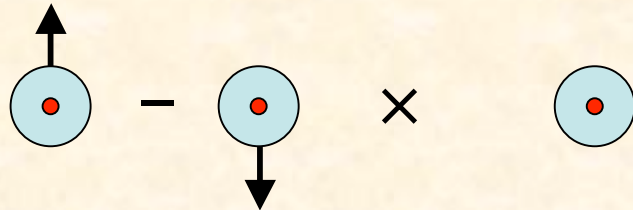


Nucleon

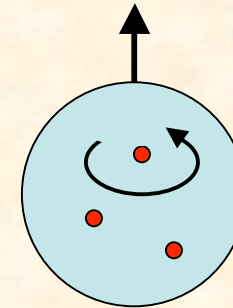


Quark

$$A_N \sim f_{1T}^\perp \cdot D_1 \quad (\text{Sivers function} \times \text{Unpolarized fragmentation})$$



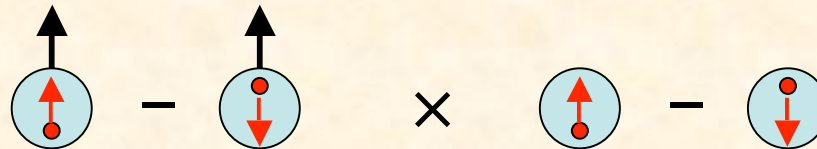
The Sivers function describes unpolarized quark in the transversely polarized nucleon.



Burkardt
@J-PARC-HS05

Probe of angular momentum

- Collins effect



$$A_N \sim \delta_T q \cdot H_1^\perp \quad (\text{Transversity} \times \text{Collins fragmentation function})$$

The transversity distribution describes transverse quark polarization in the transversely polarized nucleon.

The Collins fragmentation function describes a fragmentation of polarized quark into unpolarized hadron.

- Higher-twist

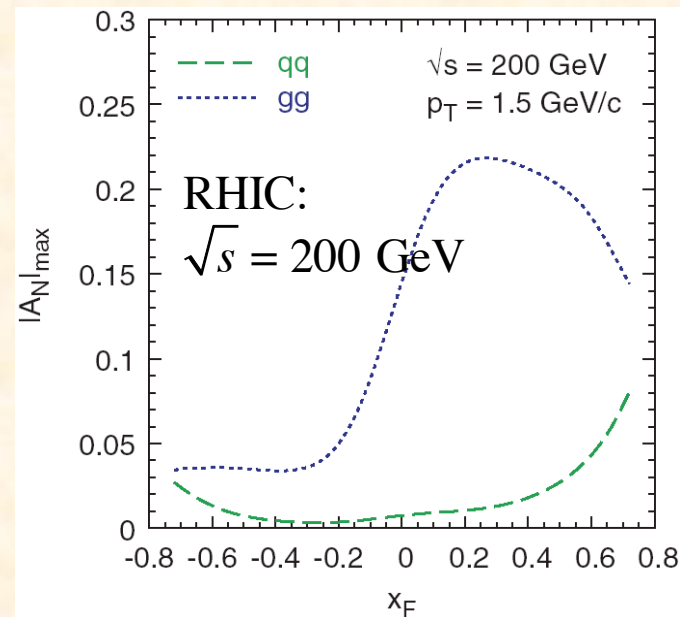
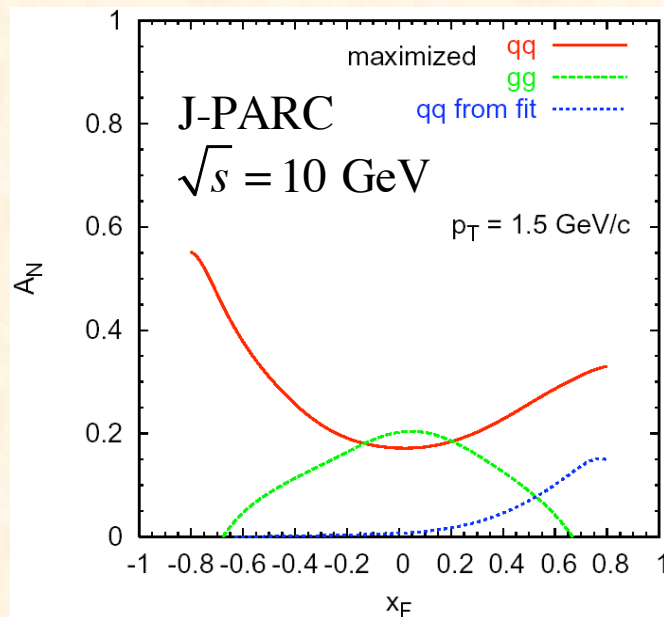
Qiu, Sterman; Koike@J-PARC-HS05

Single spin asymmetry

D-meson production

$$gg \rightarrow c\bar{c}, \quad q\bar{q} \rightarrow c\bar{c}$$

$\rightarrow c$ & \bar{c} are not polarized (no Collins mechanism)



Y. Goto
@J-PARC-HS05

U. D'Alesio
@BNL, 2005

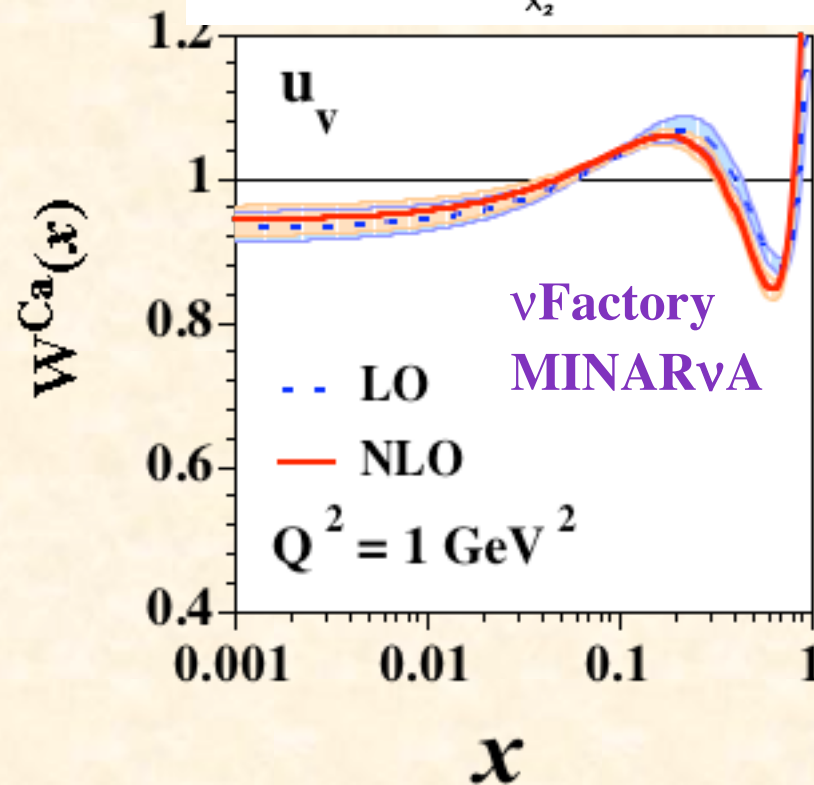
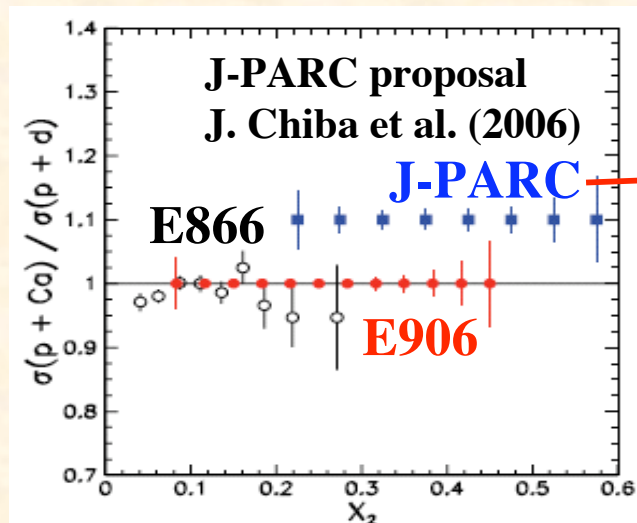
In the region $x_F < 0$

J-PARC: sensitive to quark Sivers effect

RHIC: sensitive to gluon Sivers effect

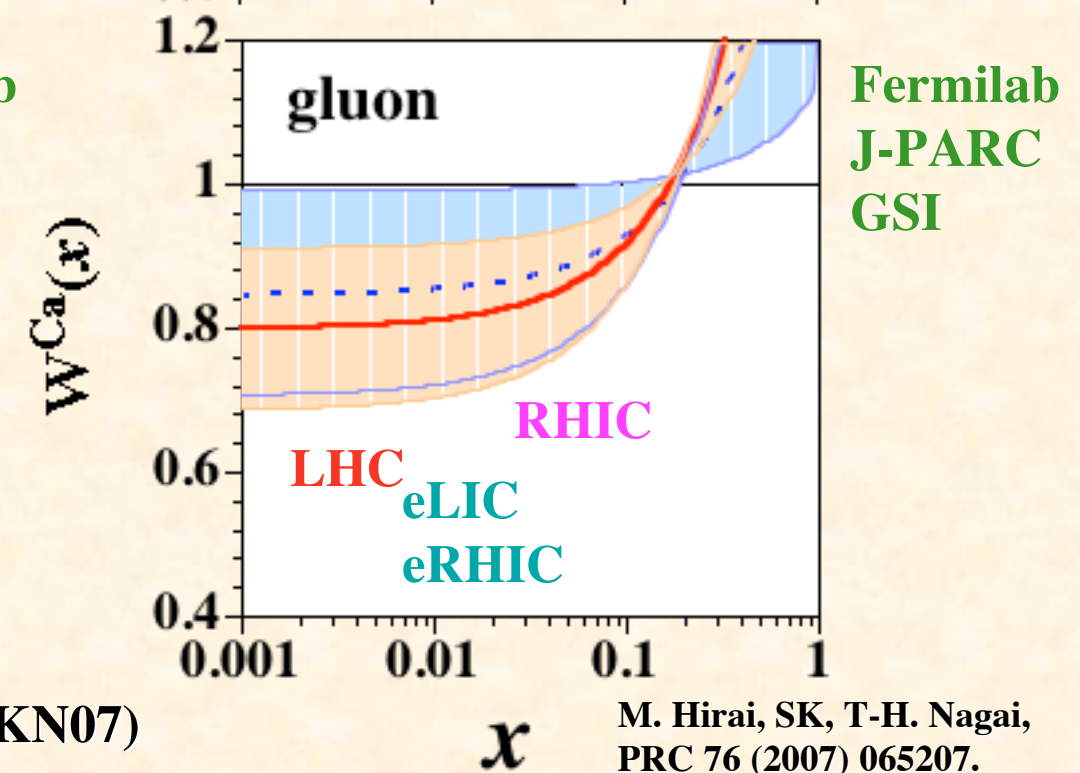
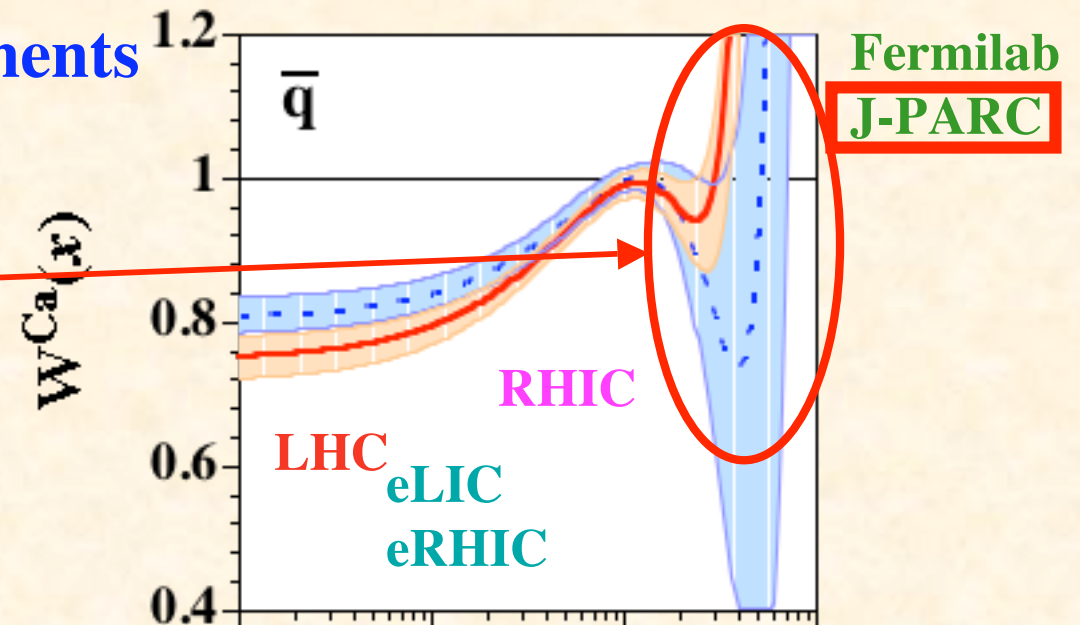
M. Anselmino *et al.*, PRD 70 (2004) 074025.

Nuclear PDFs & Future experiments



JLab

(HKN07)



M. Hirai, SK, T-H. Nagai,
PRC 76 (2007) 065207.

Short-range NN interaction

Ciofi degli Atti@J-PARC-NP07

Strikman@INPC07

E. Piasetzky et al.,
PRL97 (2006) 162504

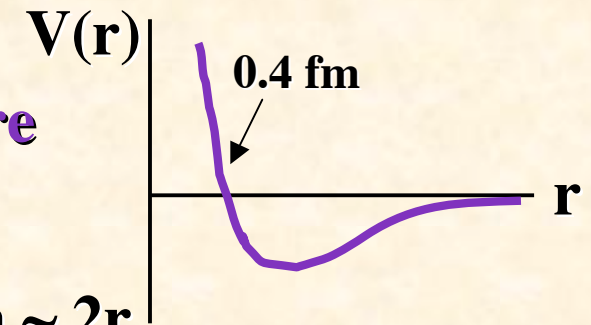
- Short-range repulsive core, Tensor force
- Quark degrees of freedom
- Cold dense nuclear matter, Neutron star

Nuclei do not collapse → Short-range repulsive core

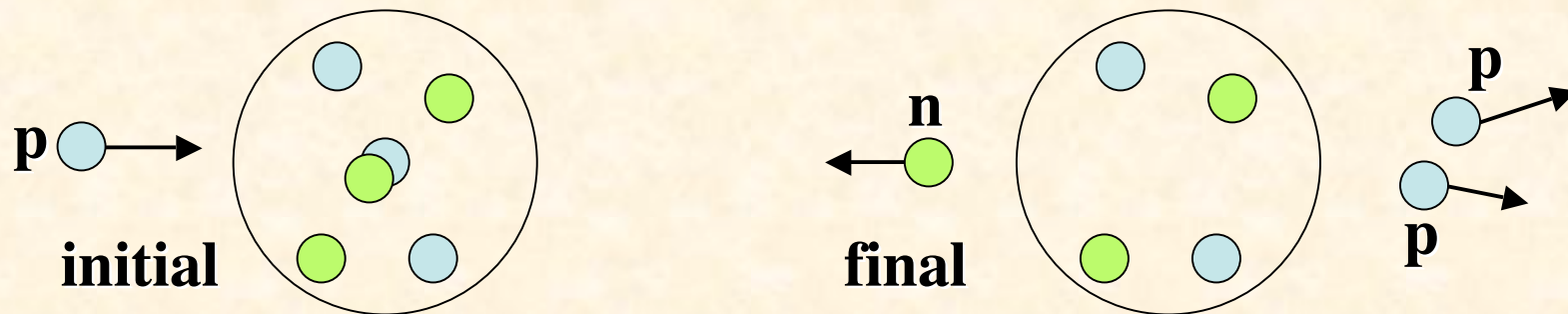
Nucleon size $r \approx 0.8$ fm

Average nucleon separation in a nucleus: $R \approx 2$ fm $\sim 2r$

→ The short-range part is important as the density becomes larger (neutron star).



$A(p, 2pN)X$ experiment for short range correlation



Color Transparency

“Probe of dynamics of elementary reactions”

At large momentum transfer, a small-size component of the hadron wave function should dominate. This small-size hadron could freely pass through nuclear medium. (Transparent)

Investigate $p A \rightarrow p p (A-1)$

Nuclear transparency: $T = \frac{\sigma_A}{A\sigma_N}$

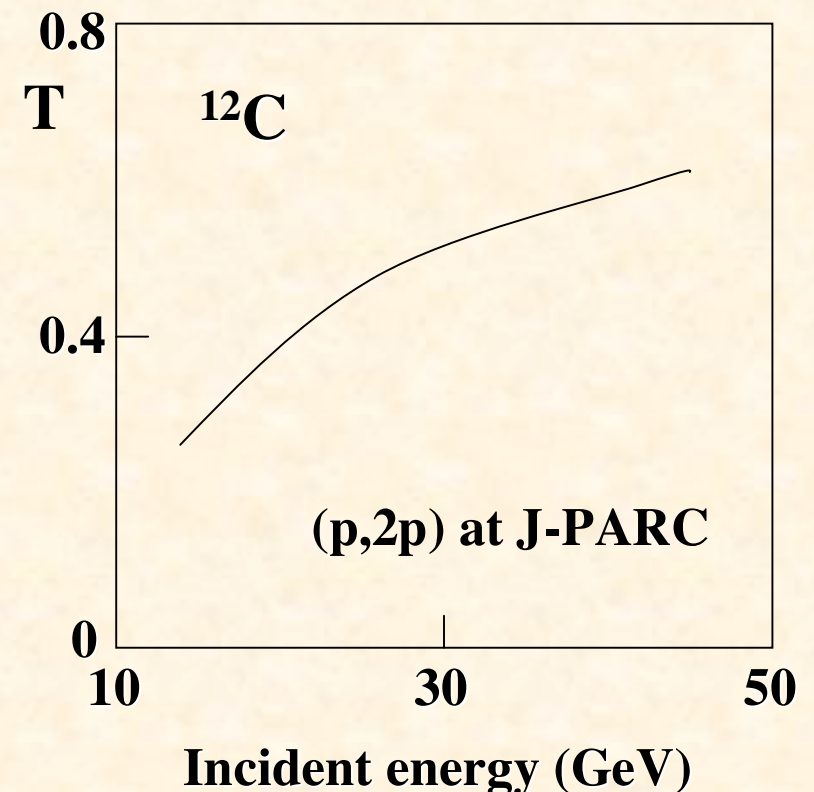
Hadron size $\sim 1 / \text{hard scale}$

Color transparency:

$T \rightarrow \text{larger, as the hard scale} \rightarrow \text{larger}$

(BNL-EVA) J. Aclander et al.,
PRC 70 (2004) 015208

Possibility at J-PARC



Possible Topics II

(closely related to my studies)

Structure of Spin-1 Hadrons

**Note: Proton-beam polarization is not needed.
Polarized deuteron target is enough at J-PARC!**

Tensor Structure in High-energy Reactions

(Note: No polarized proton beam is needed!)

L. L. Frankfurt and M. I. Strikman, NP A405 (1983) 557.

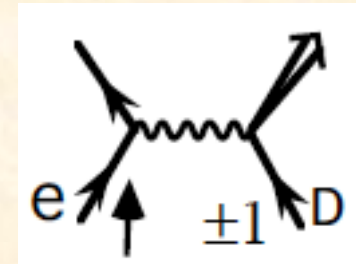
P. Hoodbhoy, R. L. Jaffe, and A. Manohar, NP B312 (1989) 571.

**Structure
Functions**
(in e scattering)

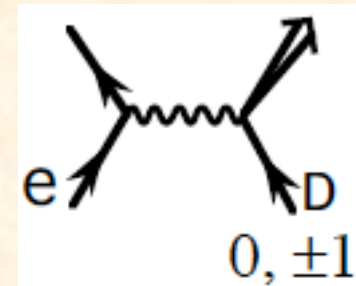
$$F_1 \propto \langle d\sigma \rangle$$



$$g_1 \propto d\sigma(\uparrow, +1) - d\sigma(\uparrow, -1)$$



$$b_1 \propto d\sigma(0) - \frac{d\sigma(+1) + d\sigma(-1)}{2}$$



**Parton
Model**

$$[q_{\uparrow}^H(x, Q^2)]$$

Unpolarized
quark distribution
in a spin-1 hadron.

$$F_1 = \frac{1}{2} \sum_i e_i^2 (q_i + \bar{q}_i)$$

$$q_i = \frac{1}{3} (q_i^{+1} + q_i^0 + q_i^{-1})$$

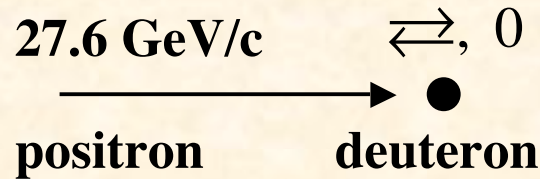
$$g_1 = \frac{1}{2} \sum_i e_i^2 (\Delta q_i + \Delta \bar{q}_i)$$

$$\Delta q_i = q_{i\uparrow}^{+1} - q_{i\downarrow}^{+1}$$

$$b_1 = \frac{1}{2} \sum_i e_i^2 (\delta q_i + \delta \bar{q}_i)$$

$$\delta q_i = q_i^0 - \frac{q_i^{+1} + q_i^{-1}}{2}$$

HERMES results on b_1



b_1 measurement in the kinematical region
 $0.01 < x < 0.45$, $0.5 \text{ GeV}^2 < Q^2 < 5 \text{ GeV}^2$

b_1 sum rule

$$\int_{0.02}^{0.85} dx b_1(x) = [1.05 \pm 0.34(\text{stat}) \pm 0.35(\text{sys})] \times 10^{-2}$$

at $Q^2 = 5 \text{ GeV}^2$

In the restricted Q^2 range $Q^2 > 1 \text{ GeV}^2$

$$\int_{0.02}^{0.85} dx b_1(x) = [0.35 \pm 0.10(\text{stat}) \pm 0.18(\text{sys})] \times 10^{-2}$$

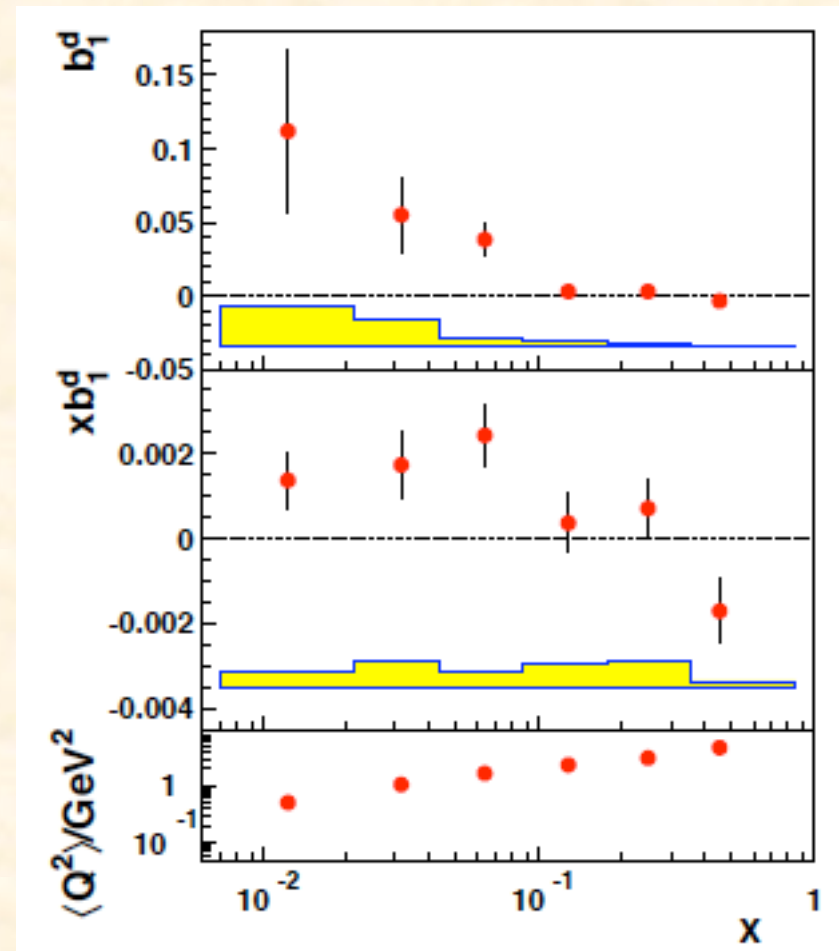
at $Q^2 = 5 \text{ GeV}^2$

F. E. Close and SK,
PRD42, 2377 (1990)

$$\int dx b_1^D(x) = \lim_{t \rightarrow 0} -\frac{5}{12} \frac{t}{M^2} F_Q(t) + \frac{1}{9} (\delta Q + \delta \bar{Q})_{\text{sea}} \Rightarrow 0$$

Electric quadrupole form factor

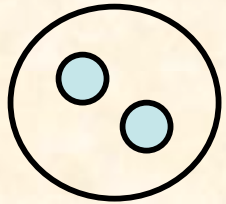
A. Airapetian *et al.*, PRL 95 (2005) 242001



Tensor Structure in Proton-Deuteron Drell-Yan

b_1 for spin-1 particles

(Note: No polarized proton beam is needed!)



only in S-wave

$$b_1 = 0$$

1st measurement of b_1 :
(HERMES) A. Airapetian et al.,
PRL 95 (2005) 242001.

Spin asymmetries

$$A_{LL} = \frac{\sum_a e_a^2 [\Delta q_a(x_A) \Delta \bar{q}_a(x_B) + \Delta \bar{q}_a(x_A) \Delta q_a(x_B)]}{\sum_a e_a^2 [q_a(x_A) \bar{q}_a(x_B) + \bar{q}_a(x_A) q_a(x_B)]}$$

$$A_{TT} = \frac{\sin^2 \theta \cos(2\phi)}{1 + \cos^2 \theta} \frac{\sum_a e_a^2 [\Delta_T q_a(x_A) \Delta_T \bar{q}_a(x_B) + \Delta_T \bar{q}_a(x_A) \Delta_T q_a(x_B)]}{\sum_a e_a^2 [q_a(x_A) \bar{q}_a(x_B) + \bar{q}_a(x_A) q_a(x_B)]}$$

$$A_{UQ_0} = \frac{\sum_a e_a^2 [q_a(x_A) \delta \bar{q}_a(x_B) + \bar{q}_a(x_A) \delta q_a(x_B)]}{\sum_a e_a^2 [q_a(x_A) \bar{q}_a(x_B) + \bar{q}_a(x_A) q_a(x_B)]}$$

$$\delta q_i = q_i^0 - \frac{q_i^{+1} + q_i^{-1}}{2}$$

Note: $\delta \neq$ transversity in my notation

Unpolarized proton
+ Tensor polarized deuteron

Unique advantage of J-PARC ($\delta \bar{q}$ measurement)

$$A_{UQ_0}(\text{large } x_F) \approx \frac{\sum_a e_a^2 q_a(x_A) \delta \bar{q}_a(x_B)}{\sum_a e_a^2 q_a(x_A) \bar{q}_a(x_B)}$$

$$\int dx b_1^D(x) = 0 + \frac{1}{9} (\delta Q + \delta \bar{Q})_{\text{sea}}$$

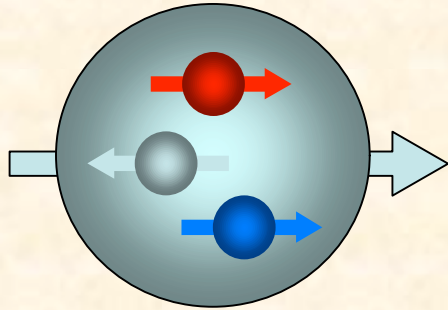
F. E. Close and SK, PRD42, 2377 (1990)

Gottfried: $\int \frac{dx}{x} [F_2^p(x) - F_2^n(x)] = \frac{1}{3} + \frac{2}{3} \int dx [\bar{u} - \bar{d}]$

Generalized Parton Distributions at Hadron Facilities (J-PARC, GSI, ...)

**Ref. Novel two-to-three hard hadronic processes and possible studies
of generalized parton distributions at hadron facilities,
SK, M. Strikman, K. Sudoh, arXiv:0905.1453 [hep-ph].**

Nucleon Spin



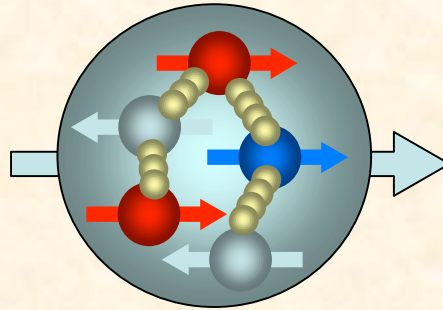
Naïve Quark Model

$$\Delta\Sigma = \Delta u_v + \Delta d_v = 1$$

Electron / muon scattering

$$\Delta\Sigma \approx 0.1 \sim 0.3$$

Almost none of nucleon spin is carried by quarks!



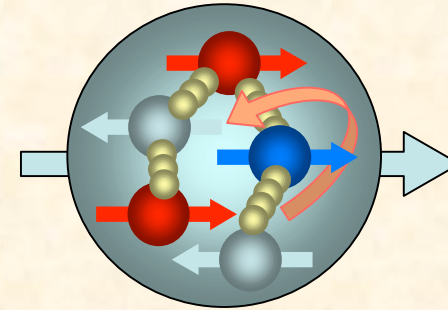
QCD

Sea-quarks and gluons?

Gluon: ΔG

Sea-quarks: Δq_{sea}

Recent data indicate ΔG is small at $x \sim 0.1$.



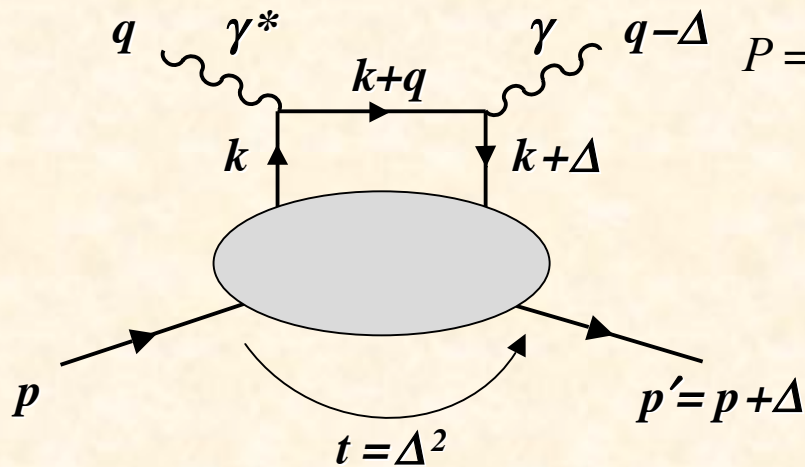
Orbital angular momenta ?

L_q, L_g

Future experiments

$$\text{Nucleon Spin: } \frac{1}{2} = \frac{1}{2} \underbrace{\left(\Delta u_v + \Delta d_v + \Delta q_{sea} \right)}_{\Delta\Sigma} + \Delta G + L_q + L_g$$

Generalized Parton Distributions (GPDs)



$$P = \frac{p + p'}{2}, \quad \Delta = p' - p$$

GPDs are defined as correlation of off-forward matrix.

Bjorken variable $x = \frac{Q^2}{2p \cdot q}$

Momentum transfer squared $t = \Delta^2$

Skewness parameter $\xi = \frac{p^+ - p'^+}{p^+ + p'^+} = -\frac{\Delta^+}{2P^+}$

$$\int \frac{dz^-}{4\pi} e^{ixP^+z^-} \langle p' | \bar{\psi}(-z/2) \gamma^+ \psi(z/2) | p \rangle \Big|_{z^+=0, \vec{z}_\perp=0} = \frac{1}{2P^+} \left[H(x, \xi, t) \bar{u}(p') \gamma^+ u(p) + E(x, \xi, t) \bar{u}(p') \frac{i\sigma^{+\alpha} \Delta_\alpha}{2M} u(p) \right]$$

$$\int \frac{dz^-}{4\pi} e^{ixP^+z^-} \langle p' | \bar{\psi}(-z/2) \gamma^+ \gamma_5 \psi(z/2) | p \rangle \Big|_{z^+=0, \vec{z}_\perp=0} = \frac{1}{2P^+} \left[\tilde{H}(x, \xi, t) \bar{u}(p') \gamma^+ \gamma_5 u(p) + \tilde{E}(x, \xi, t) \bar{u}(p') \frac{\gamma_5 \Delta^+}{2M} u(p) \right]$$

Forward limit: PDFs

$$H(x, \xi, t) \Big|_{\xi=t=0} = f(x), \quad \tilde{H}(x, \xi, t) \Big|_{\xi=t=0} = \Delta f(x)$$

First moments: Form factors

There is no analog in E and \tilde{E} .

Dirac and Pauli form factors F_1, F_2

$$\int dx H(x, \xi, t) = F_1(t), \quad \int dx E(x, \xi, t) = F_2(t)$$

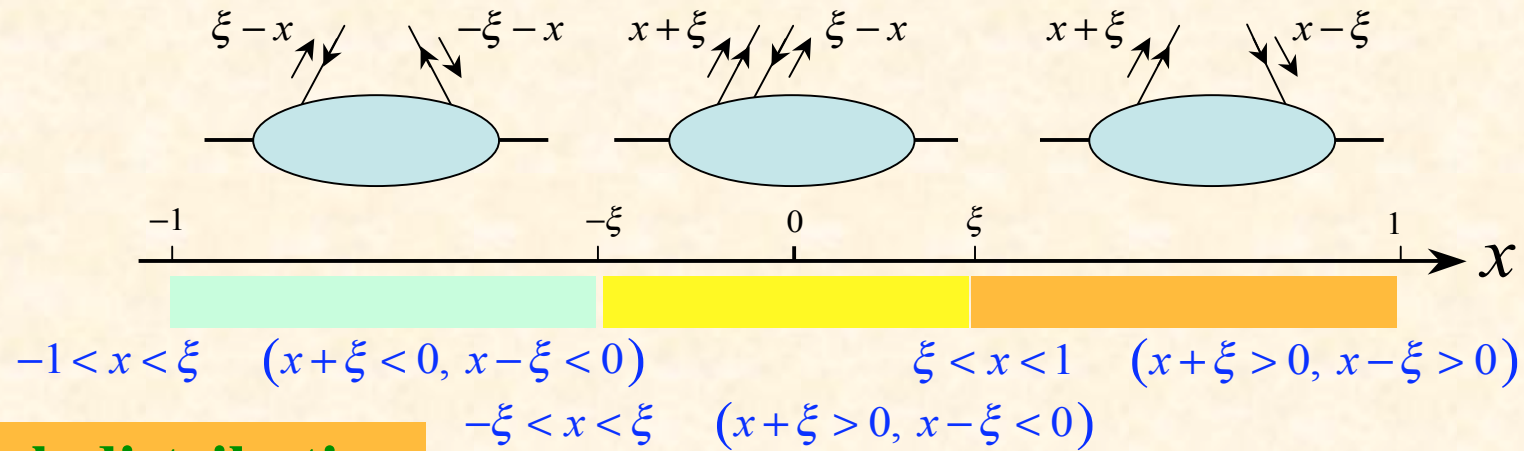
Axial and Pseudoscalar form factors G_A, G_P

$$\int dx \tilde{H}(x, \xi, t) = G_A(t), \quad \int dx \tilde{E}(x, \xi, t) = G_P(t)$$

Second moments: Angular momenta

Sum rule: $J_q = \frac{1}{2} \int dx x [H_q(x, \xi, t=0) + E_q(x, \xi, t=0)], \quad J_q = \frac{1}{2} \Delta q + L_q$

GPDs in different x regions and GPDs at J-PARC



Quark distribution

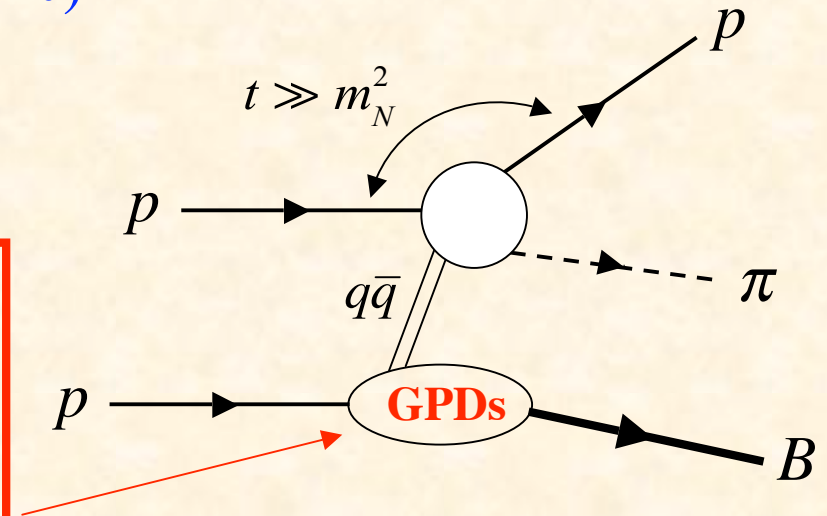
Emission of quark with momentum fraction $x+\xi$
 Absorption of quark with momentum fraction $x-\xi$

Meson distribution amplitude

Emission of quark with momentum fraction $x+\xi$
 Emission of antiquark with momentum fraction $\xi-x$

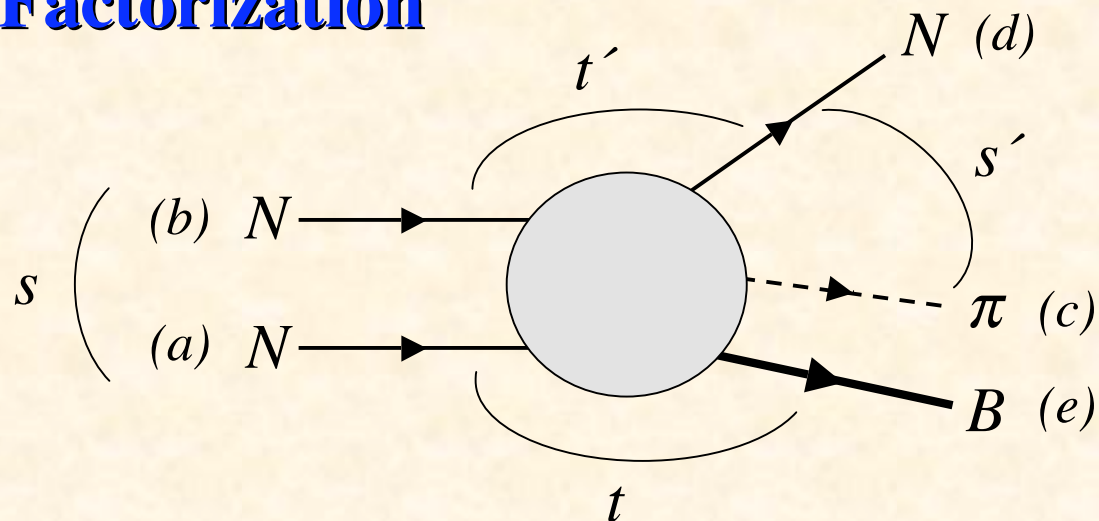
Antiquark distribution

Emission of antiquark with momentum fraction $\xi-x$
 Absorption of antiquark with momentum fraction $-x-\xi$



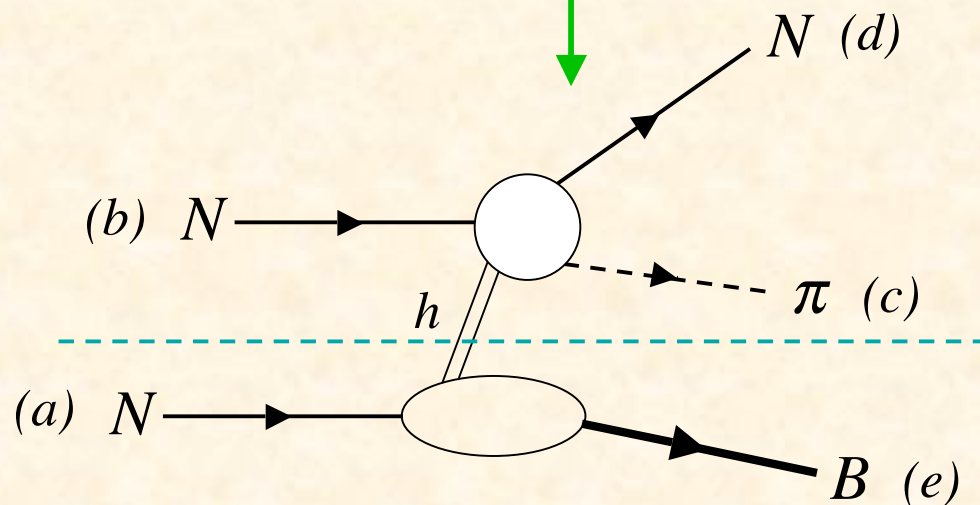
GPDs at J-PARC:
 SK, M. Strikman, K. Sudoh,
 arXiv:0905.1453

Factorization

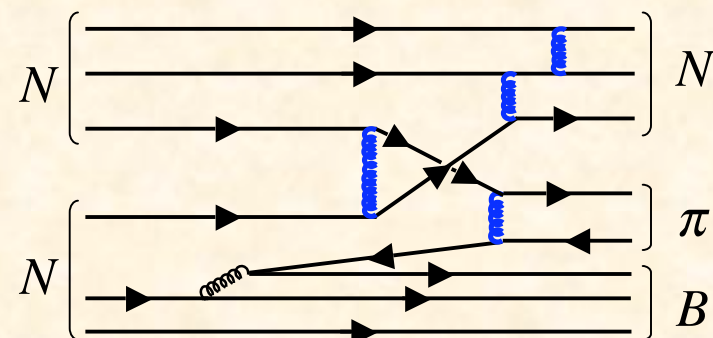


Consider a hard reaction with

$$|s'|, |t'|, |u'| \gg M_N^2, |t| \ll M_N^2$$

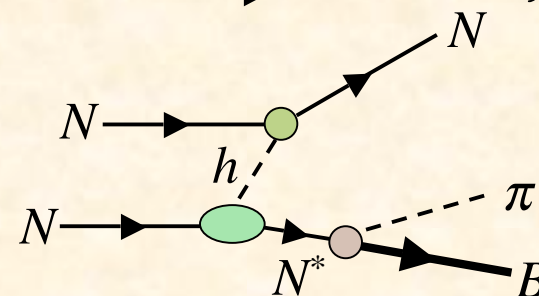
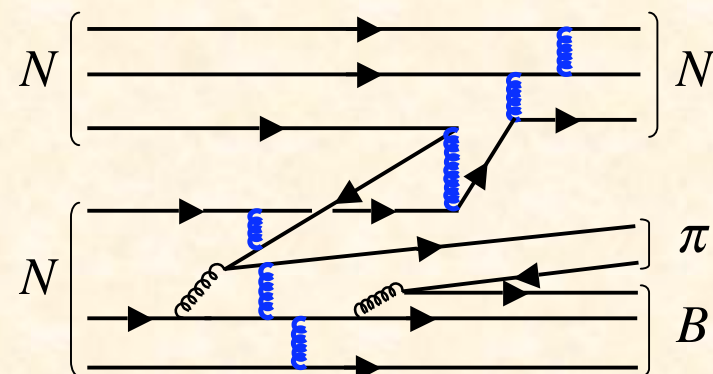


Typical leading process

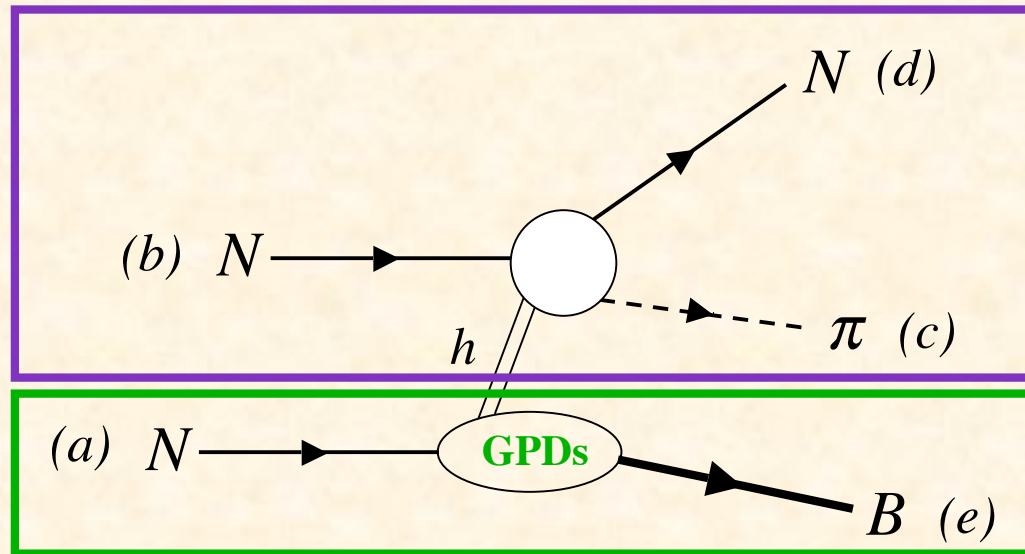


Typical sub-leading process
(cannot be factorized)

More hard gluon exchanges
→ suppressed

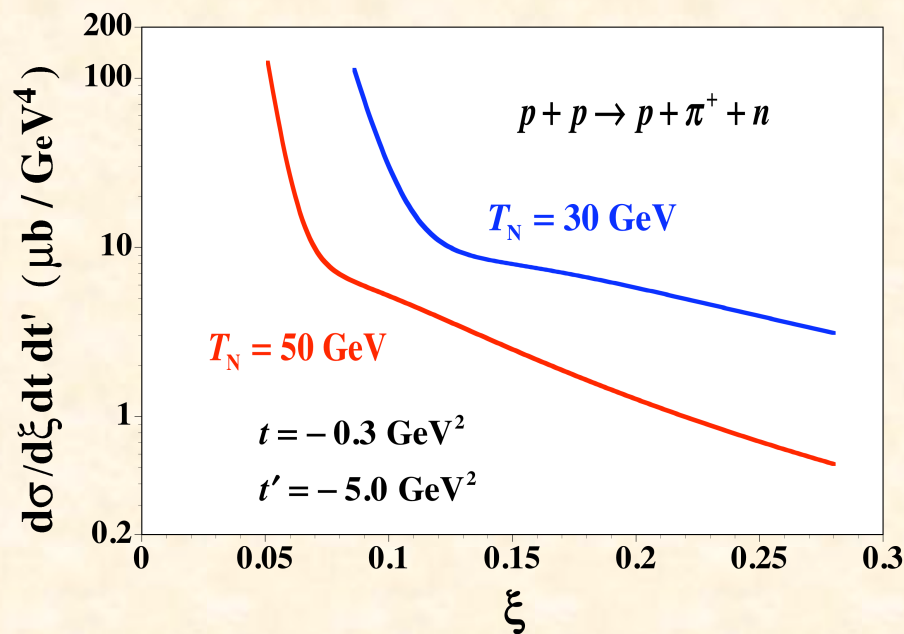


Cross section estimates



$\frac{d\sigma(s', t')}{dt'}$ so as to explain
AGS experimental data on
 $\pi + p \rightarrow \pi + p, \pi + p \rightarrow \rho + p$

This part is expressed by GPDs.



At this stage, our numerical results are
for rough order of magnitude estimates
on cross sections by assuming π - and ρ -like
intermediate states.

Still premature stage, need more works!

Exotic hadron search by fragmentation functions

**Refs. (1) Scalar mesons in ϕ radiative decay,
F. E. Close, N. Isgur, S. Kumano,
Nucl. Phys. B389 (1993) 513.**

**(2) Proposal for exotic-hadron search
by fragmentation functions,
M. Hirai, S. Kumano, M. Oka, K. Sudoh,
Phys. Rev. D77 (2008) 017504.**

Recent progress in exotic hadrons

$q\bar{q}$ Meson
 q^3 Baryon

$q^2\bar{q}^2$ Tetraquark
 $q^4\bar{q}$ Pentaquark
 q^6 Dibaryon

...
 $q^{10}\bar{q}$ e.g. Strange
 tribaryon

...
 gg Glueball
 ...

(Japanese ?) Exotics

- $\Theta^+(1540)?$: LEPS

$uudd\bar{s}$?

Pentaquark?

- $S^0(3115)$, $S^+(3140)$: KEK-PS

$K^- pnn$
 $K^- ppn$?

Strange tribaryons?

- $X(3872)$, $Y(3940)$: Belle

$c\bar{c}$
 $D^0(c\bar{u})\bar{D}^0(\bar{c}u)$
 $D^+(c\bar{d})D^-(\bar{c}d)$?

Tetraquark, $D\bar{D}$ molecule

- $D_{sJ}(2317)$, $D_{sJ}(2460)$: BaBar, CLEO, Belle

Tetraquark, DK molecule

$c\bar{s}$
 $D^0(c\bar{u})K^+(u\bar{s})$
 $D^+(c\bar{d})K^0(d\bar{s})$?

- $Z(4430)$: Belle

Tetraquark, ...

Note: $Z(4430) \neq q\bar{q}$

$c\bar{c}u\bar{d}$, D molecule?

Criteria for determining f_0 structure by its fragmentation functions

M. Hirai, SK, M. Oka, K. Sudoh,
PRD 77 (2008) 017504.

Possible configurations of $f_0(980)$

- (1) ordinary u, d - meson $\frac{1}{\sqrt{2}}(u\bar{u} + d\bar{d})$
- (2) strange meson, $s\bar{s}$
- (3) tetraquark ($K\bar{K}$), $\frac{1}{\sqrt{2}}(u\bar{u}s\bar{s} + d\bar{d}s\bar{s})$
- (4) glueball gg

Contradicts with experimental widths

$$\begin{aligned}\Gamma_{\text{theo}}(f_0 \rightarrow \pi\pi) &= 500 - 1000 \text{ MeV} \\ &\gg \Gamma_{\text{exp}} = 40 - 100 \text{ MeV} \\ \Gamma_{\text{theo}}(f_0 \rightarrow \gamma\gamma) &= 1.3 - 1.8 \text{ keV} \\ &\gg \Gamma_{\text{exp}} = 0.205 \text{ keV}\end{aligned}$$

Contradicts with lattice-QCD estimate

$$\begin{aligned}m_{\text{lattice}}(f_0) &= 1600 \text{ MeV} \\ &\gg m_{\text{exp}} = 980 \text{ MeV}\end{aligned}$$

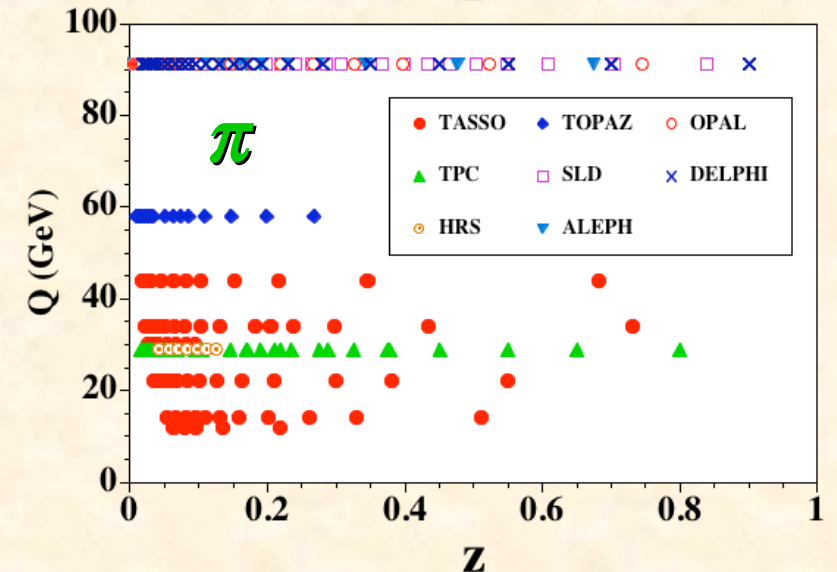
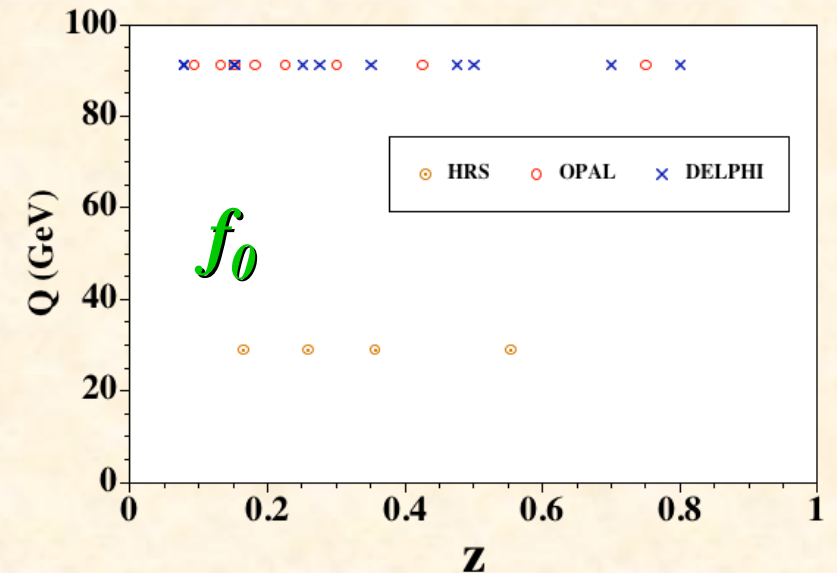
Discuss 2nd moments and functional forms (peak positions) of the fragmentation functions for f_0 by assuming the above configurations, (1), (2), (3), and (4).

Experimental data for f_0

Total number of data: **only 23**

Exp. collaboration	\sqrt{s} (GeV)	# of data
HRS	29	4
OPAL	91.2	8
DELPHI	91.2	11

One could foresee the difficulty
in getting reliable FFs for f_0
at this stage.



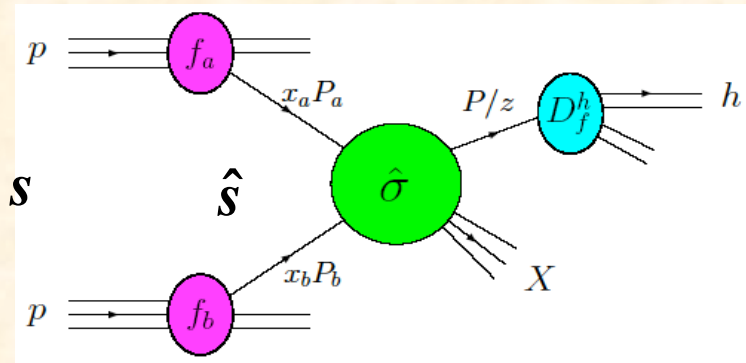
pion Total number of data: **264**

Exp. collaboration	\sqrt{s} (GeV)	# of data
TASSO	12,14,22,30,34,44	29
TCP	29	18
HRS	29	2
TOPAZ	58	4
SLD	91.2	29
SLD [light quark]		29
SLD [c quark]		29
SLD [b quark]		29
ALEPH	91.2	22
OPAL	91.2	22
DELPHI	91.2	17
DELPHI [light quark]		17
DELPHI [b quark]		17

Fragmentation functions at J-PARC

Gluon and light-quark
fragmentation functions
have large uncertainties.

Large differences between
the functions of various
analysis groups.



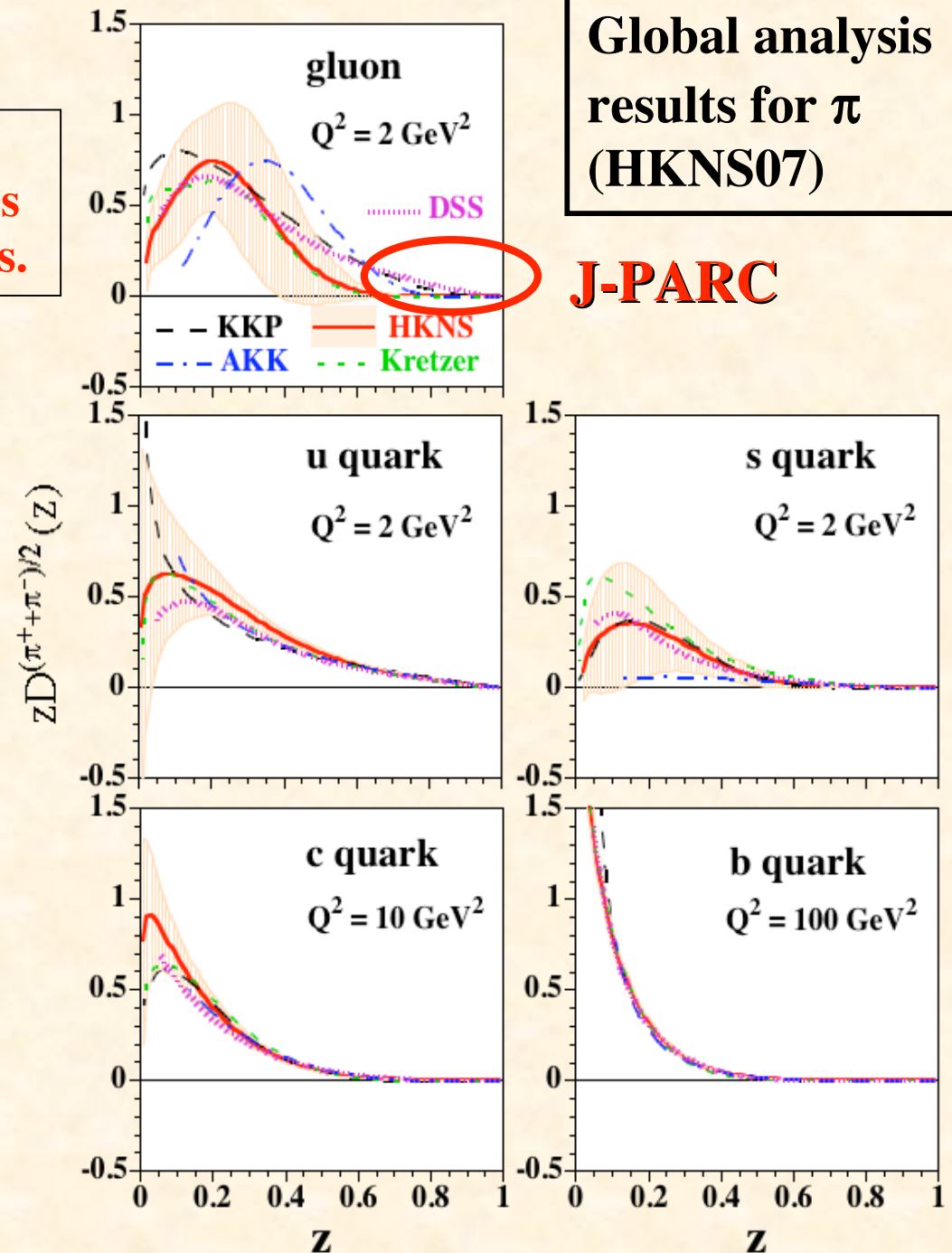
$$\hat{s} = x_a x_b s \sim (0.4)^2 (10 \text{ GeV})^2$$

$$\sqrt{\hat{s}} = 0.4 \cdot 10 = 4 \text{ GeV}$$

$$z \sim \frac{p_T}{\sqrt{\hat{s}}/2} = \frac{p_T}{2} \sim 1 \text{ (large } z\text{)}$$

Global analysis
results for π
(HKNS07)

J-PARC



Final comments and Summary

Issues for high-energy hadron project at J-PARC

- 30 GeV proton beam

- **Drell-Yan** may not be an appropriate experiment because data above the J/ψ mass are usually taken for PDF studies.
If data below the J/ψ can be used, it is a feasible experiment in the beginning stage of J-PARC.

- **Charmed-meson** production processes should be possible.

Two theoretical issues:

- (1) pQCD corrections are larger than the ones of Drell-Yan, and they may not be theoretically under control.
- (2) Production mechanisms may not be well established.

- **Better ideas** appropriate for the 30 GeV beam.

Elastic spin asymmetries, GPDs, ...

- 50 GeV proton beam

- **Drell-Yan** is possible. (First P906 at Fermilab in ~2011-2013, then at J-PARC.)
pQCD corrections are theoretically under control.

- Many other *hard* processes ... but the issue is 30→50 GeV cost.

Summary

There are interesting topics which could be investigated by using the primary 30-50 GeV proton beam.

- **Structure functions, spin physics, parton-energy loss, ...**
- **Charm nuclear physics**
- **Large- x part of parton distribution functions**
- **Perturbative QCD**
- **Tensor structure at high energy**
- **Exotic hadron search by fragmentations**
- **Transition from quark to hadron degrees of freedom**
- **Short-range nuclear force**
- **Color transparency**
- **Generalized parton distributions**

Need

- **theoretical ideas and experimental feasibility studies**
- **attractive slogan(s) and public relations**

The End

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